



Session 4: Detectors – O20

Determination of Radiation Hardness of Silicon Diodes.



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ICNMTA 2018
Guildford, Surrey, UK



What



Object of the research

Study of the radiation hardness of a commercially available silicon photo-diode commonly used as a nuclear detector

Tool

Focused MeV ion beams

to induce the damage

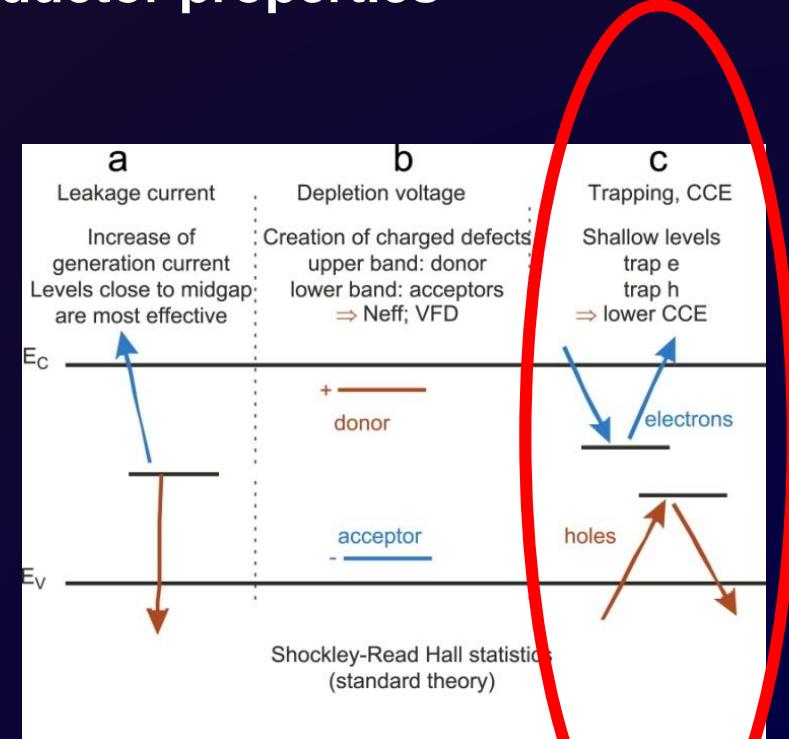
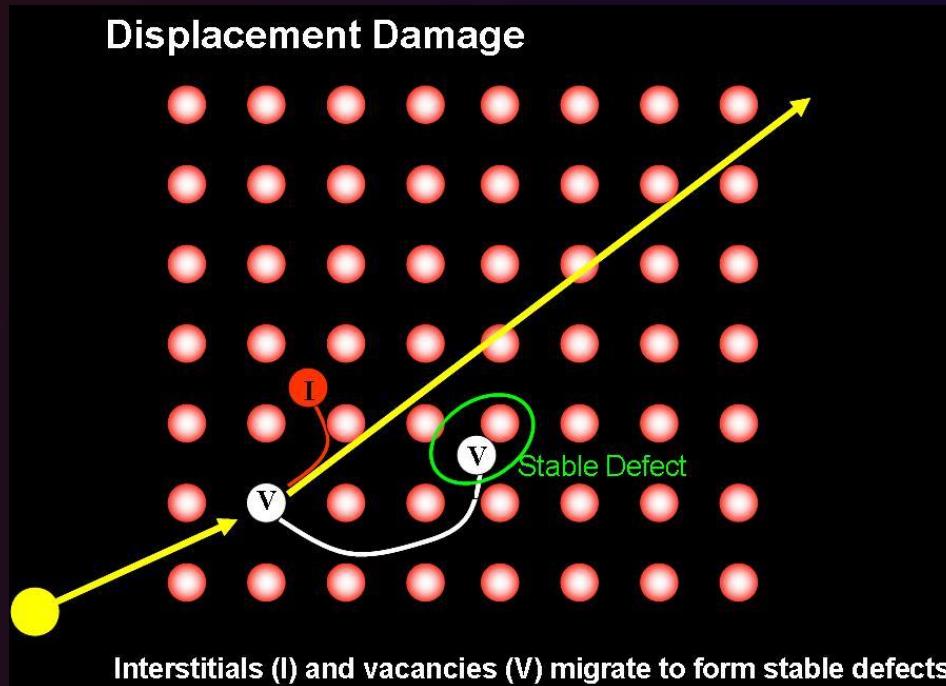
to probe the damage





Radiation damage is the general alteration of the operational properties of a semiconductor devices induced by ionizing radiation

- Displacements. Dislocations of atoms from their normal sites in the lattice, producing less ordered structures, with long term effects on semiconductor properties





WHY



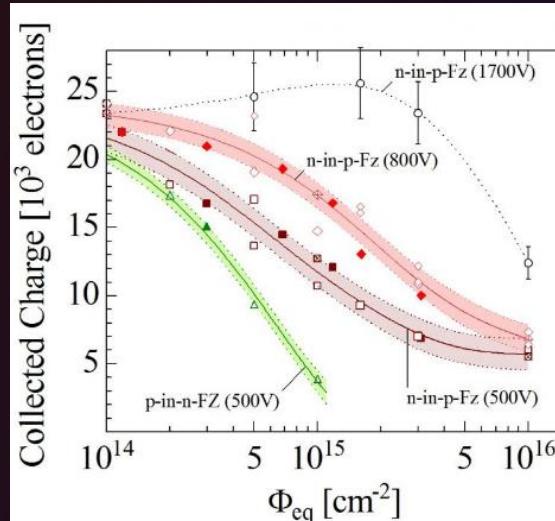
Nuclear Instruments and Methods in Physics Research A 426 (1999) 1–15

**NUCLEAR
INSTRUMENTS
& METHODS
IN PHYSICS
RESEARCH**
Section A



Radiation hardness of silicon detectors – a challenge from high-energy physics

G. Lindström*, M. Moll, E. Fretwurst



FZ Silicon Strip Sensors

- n-in-p (FZ), 300µm, 500V, 23GeV p
- n-in-p (FZ), 300µm, 500V, neutrons
- △ n-in-p (FZ), 300µm, 500V, 26MeV p
- ◆ n-in-p (FZ), 300µm, 800V, 23GeV p
- ◇ n-in-p (FZ), 300µm, 800V, neutrons
- ◆ n-in-p (FZ), 300µm, 800V, 26MeV p
- n-in-p (FZ), 300µm, 1700V, neutrons
- ▲ p-in-n (FZ), 300µm, 500V, 23GeV p
- △ p-in-n (FZ), 300µm, 500V, neutrons

RD50 - Radiation hard semiconductor devices for very high luminosity colliders



National Aeronautics and Space Administration



APPLICATIONS

The technology has several potential applications:

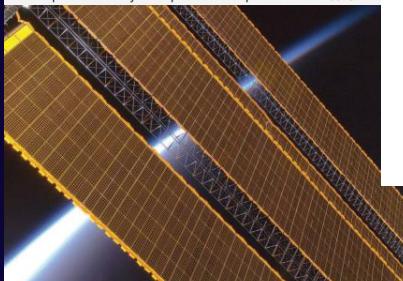
- Solar cell monitoring for manned and unmanned spacecraft
- Diagnostics for terrestrial solar power generation systems

Instrumentation

Method and Apparatus for In Situ Monitoring of Solar Cells

A novel approach to solar cell monitoring

NASA's Glenn Research Center has developed a method and apparatus for in situ health monitoring of solar cells. The innovation is a novel approach to solar cell monitoring, as it is radiation-hard, consumes few system resources, and uses commercially available components. The system operates at temperatures from -55°C to



PUBLICATIONS

Patent No: 8,159,238; 9,419,558

Patent Pending



WHY

it is relevant to the ICNMTA community



10 contributions mentioning STIM
6 contributions mentioning IBIC

Biomedical Applications
Facilities & Techniques
Detectors
Quantum Devices



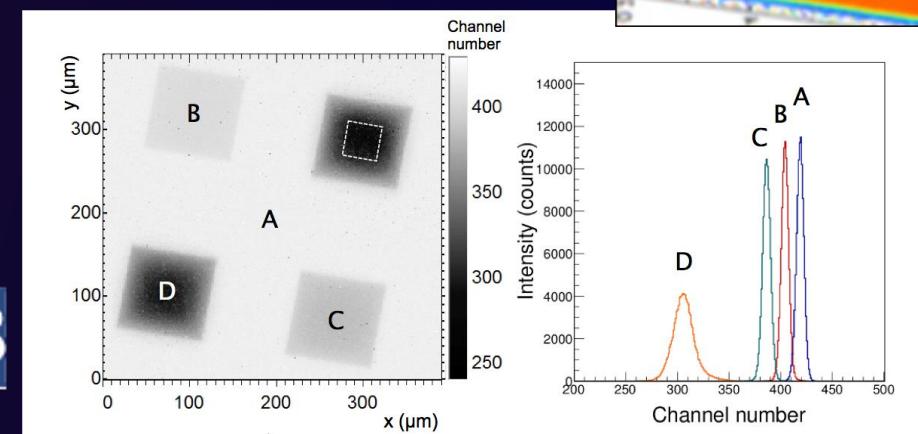
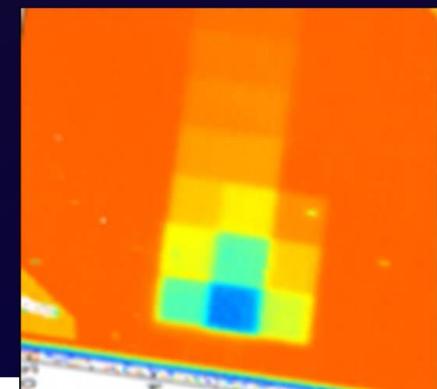
8 contributions mentioning STIM



13 contributions mentioning STIM



15 contributions mentioning STIM



Credit: Milko Jaksic





Characterization of radiation induced damage:



Device characteristic

after irradiation

before irradiation

$$\eta = \frac{Y}{Y_0} = 1 - K \cdot \Phi = 1 - K_{ed} \cdot D_d$$

Particle Fluence

Equivalent damage factor

Displacement dose

A large yellow L-shaped arrow on the left points from "before irradiation" to "after irradiation". Red arrows point from "Particle Fluence", "Equivalent damage factor", and "Displacement dose" to their respective terms in the equation.

First order: proportionality, independent of the particle, between the damage factor and the particle NIEL

NIEL approach:

measurement of K_{ed} only for one particle (at one specific energy)

K_{ed} can be estimated for all the particles and energies



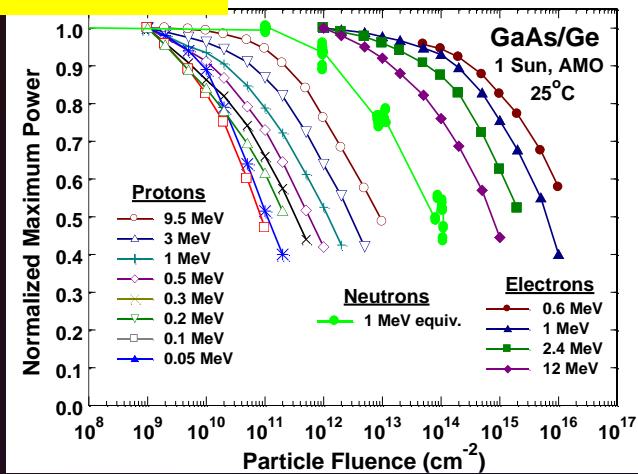


US Naval Research Laboratory (NRL)

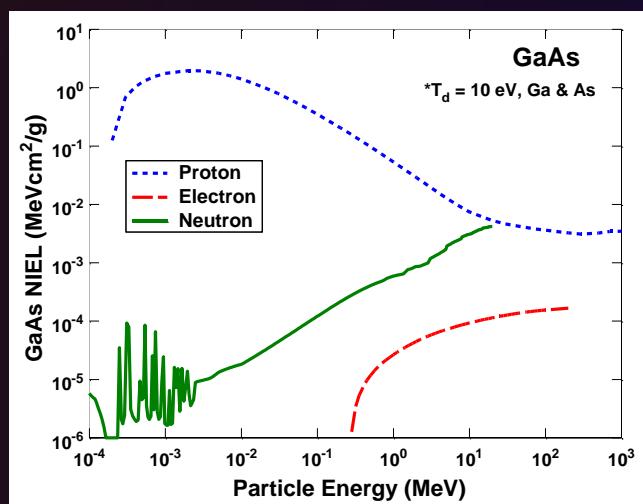
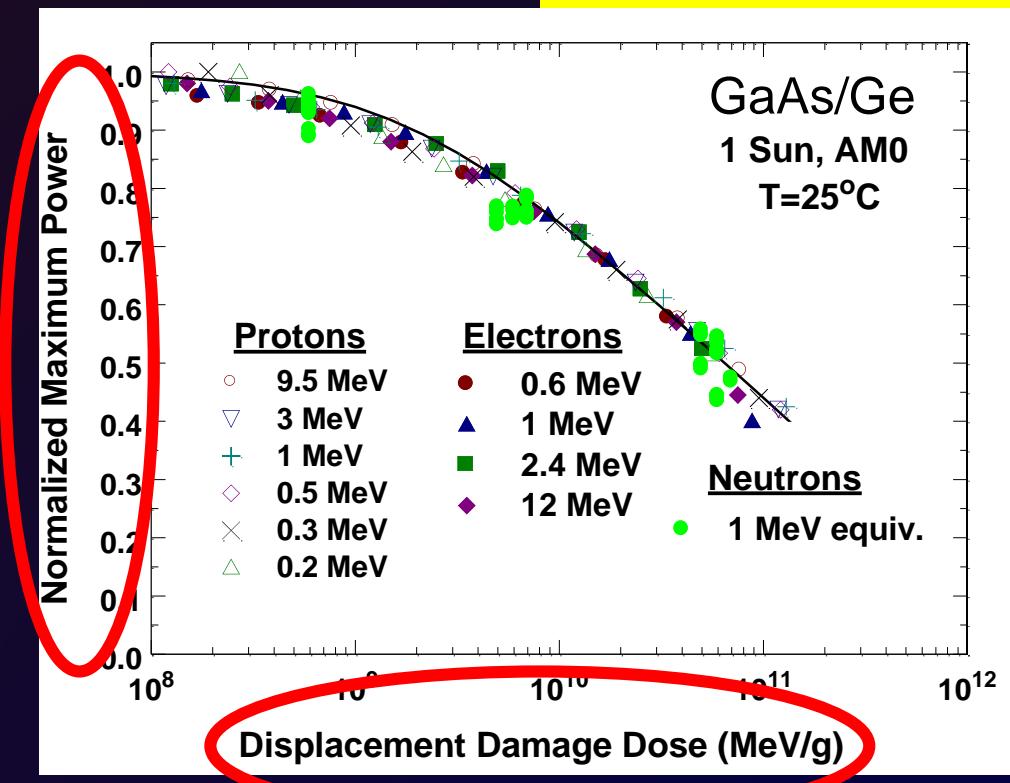
Displacement Damage Dose Method



Measured Data



Characteristic Curve



- Characteristic curve is independent of particle
- Calculated NIEL gives energy dependence of damage coefficients

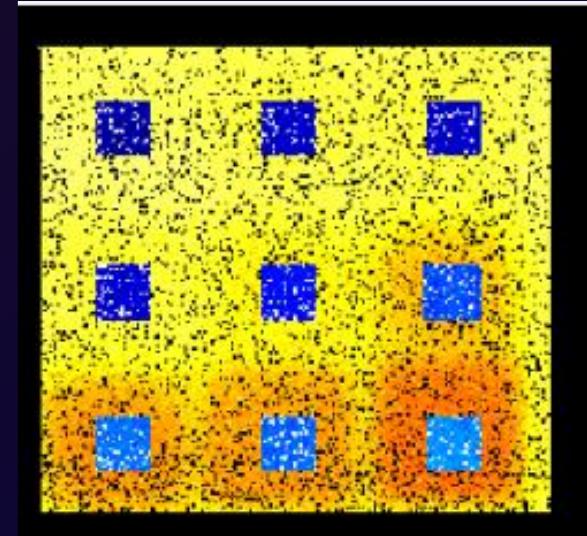
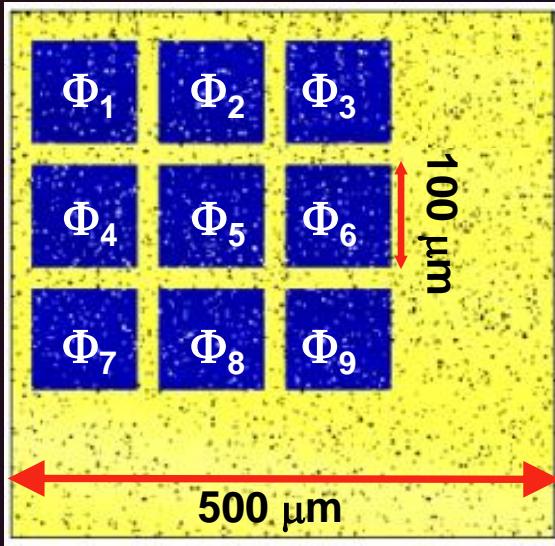


Physical Observable Charge Collection Efficiency

Focused MeV Ion beams

to induce the damage

to probe the damage



PIB
probing ion beam





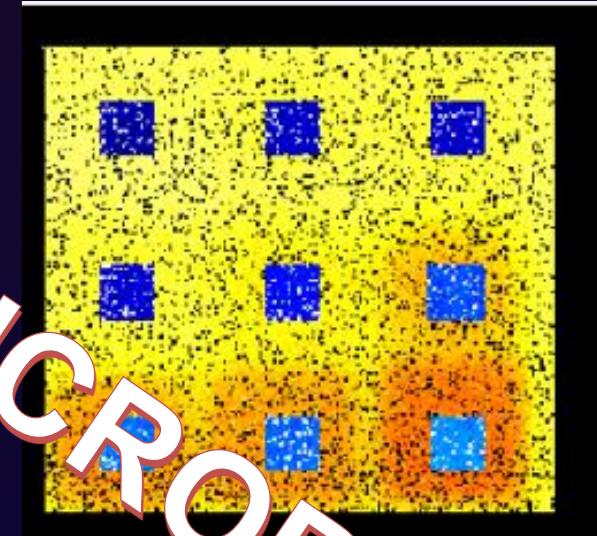
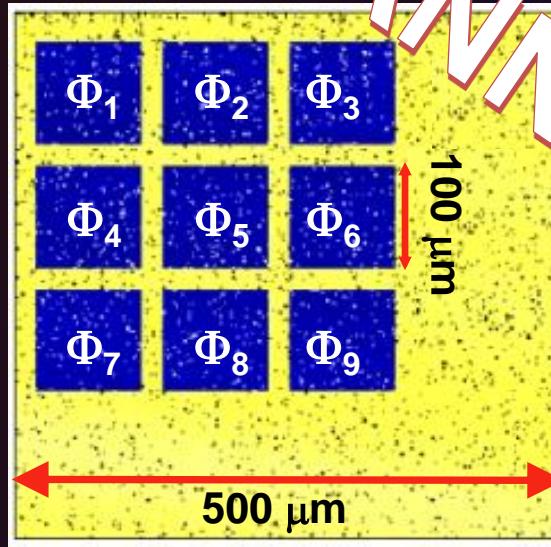
Physical Observable

Charge Collection Efficiency

Focused MeV Ion beams

to induce the damage

to probe the damage



SCANNING MICROBEAM

PIB
probing ion beam



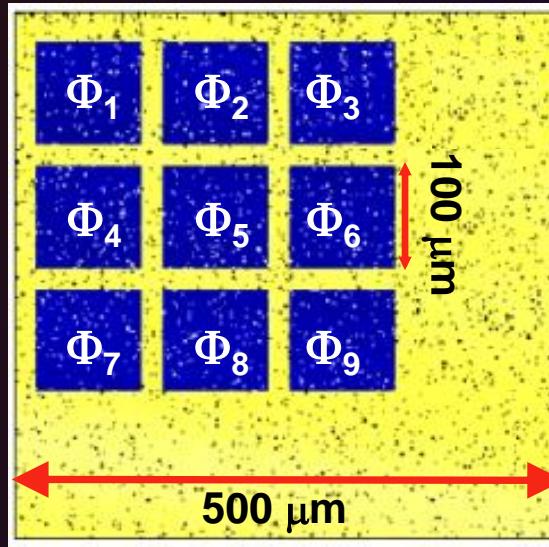


Physical Observable Charge Collection Efficiency



Focused MeV Ion beams

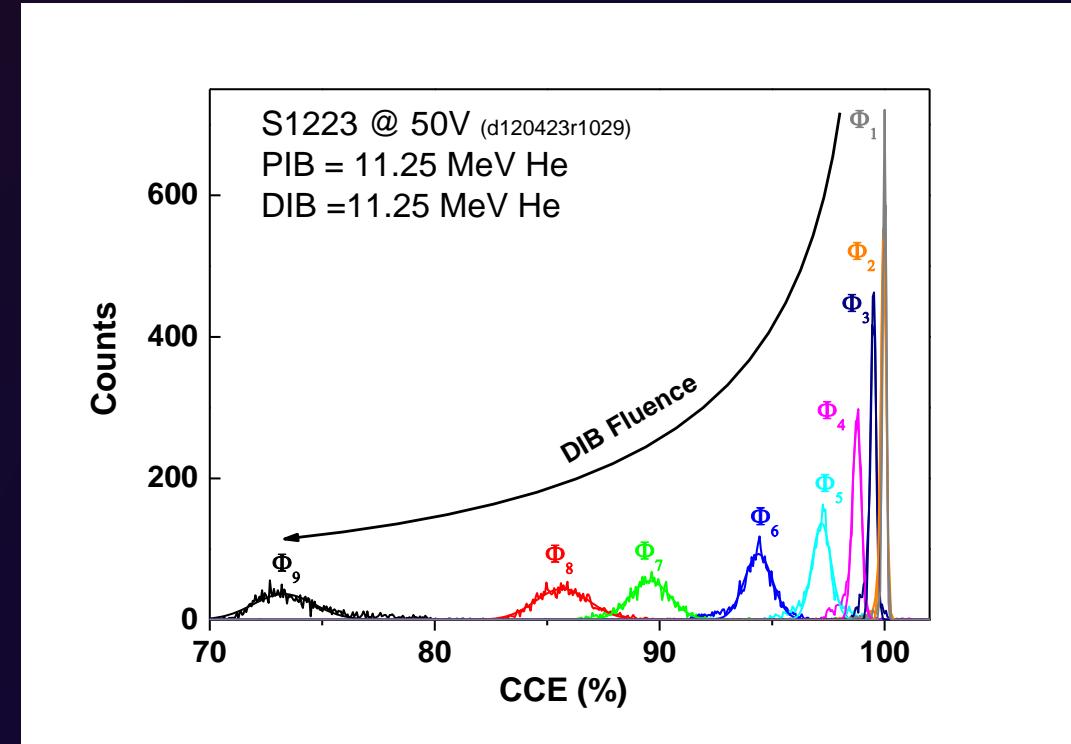
to induce the damage



DIB

damaging ion beam

to probe the damage



PIB
probing ion beam

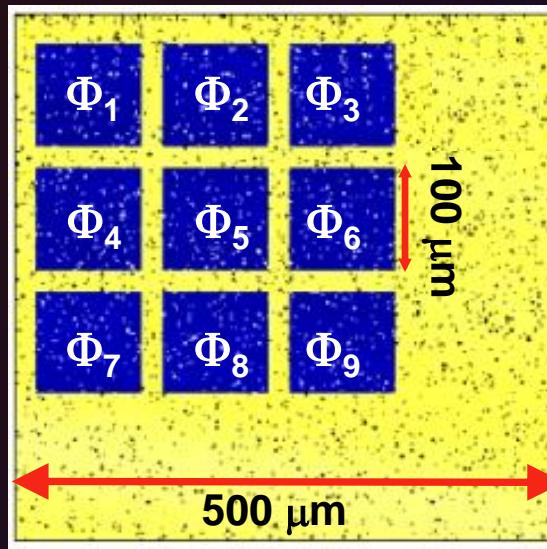




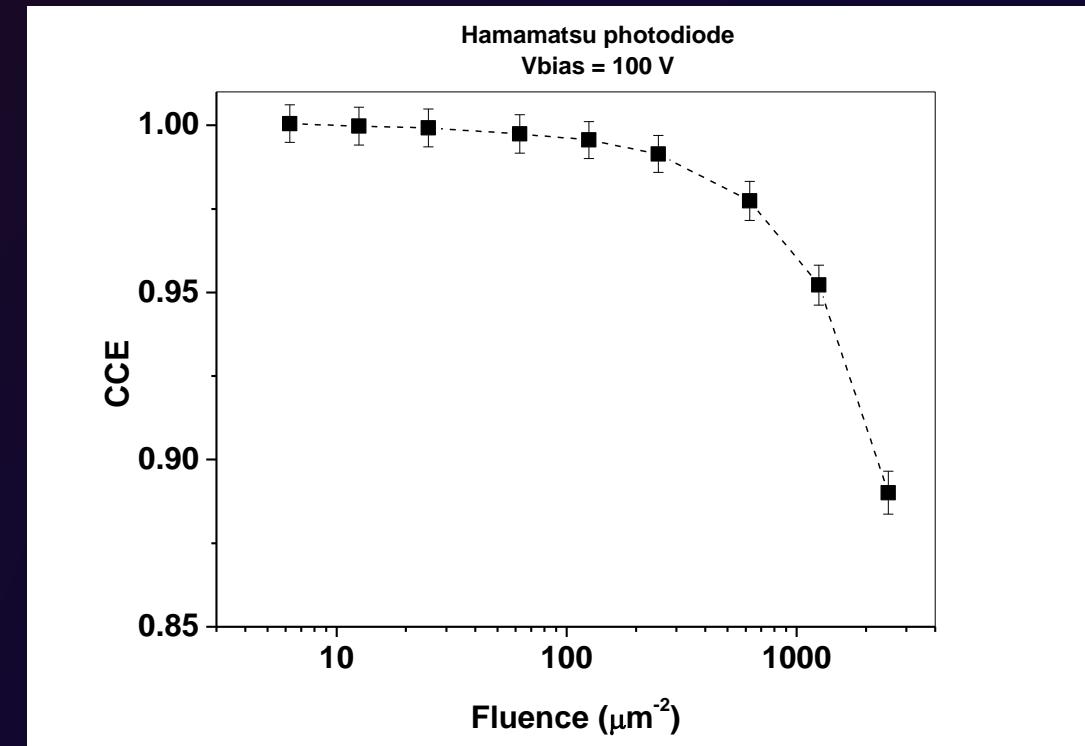
Physical Observable Charge Collection Efficiency



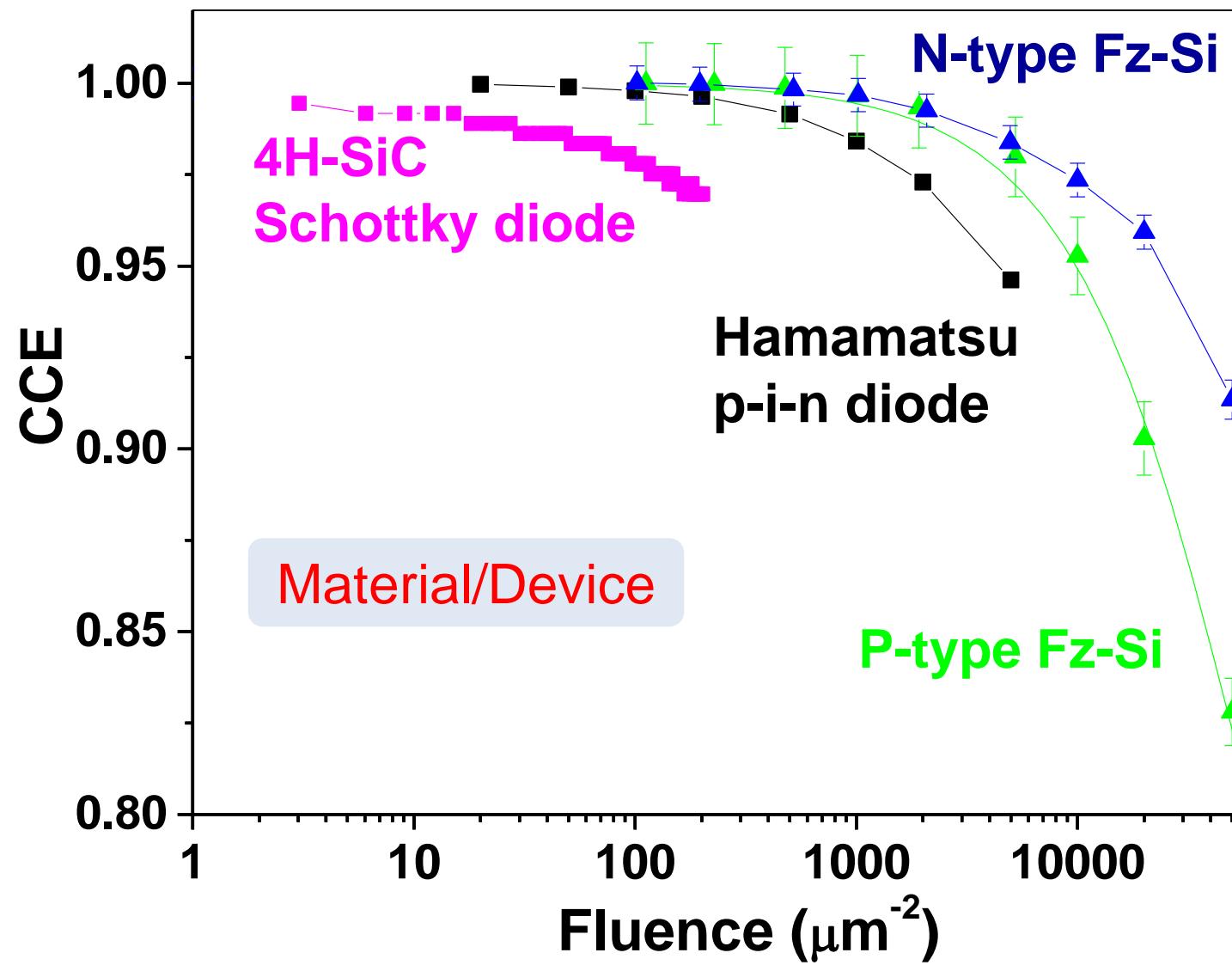
Focused MeV ion beams
to induce the damage to probe the damage



DIB = 2.15 MeV Li
damaging ion beam

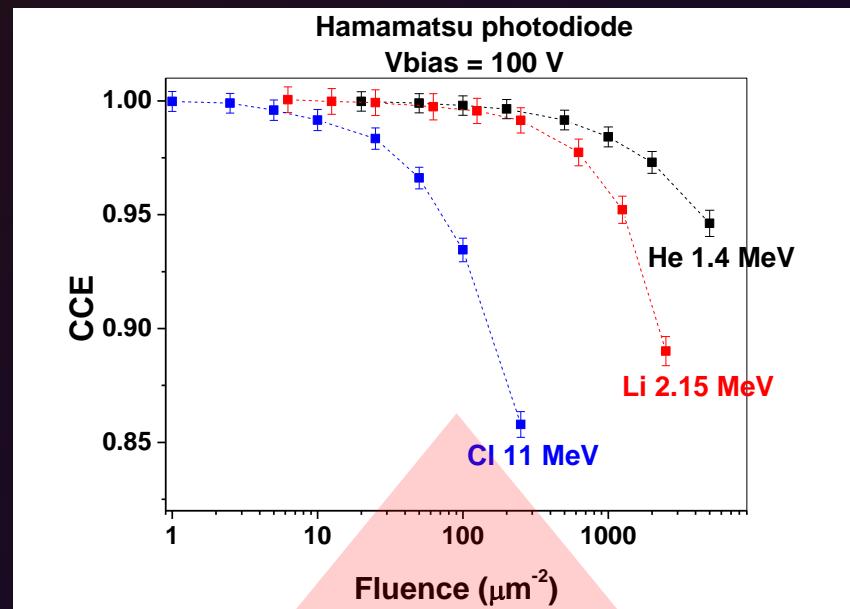


PIB = 1.4 MeV He
probing ion beam





Damaging Ion Mass/Energy Fluence

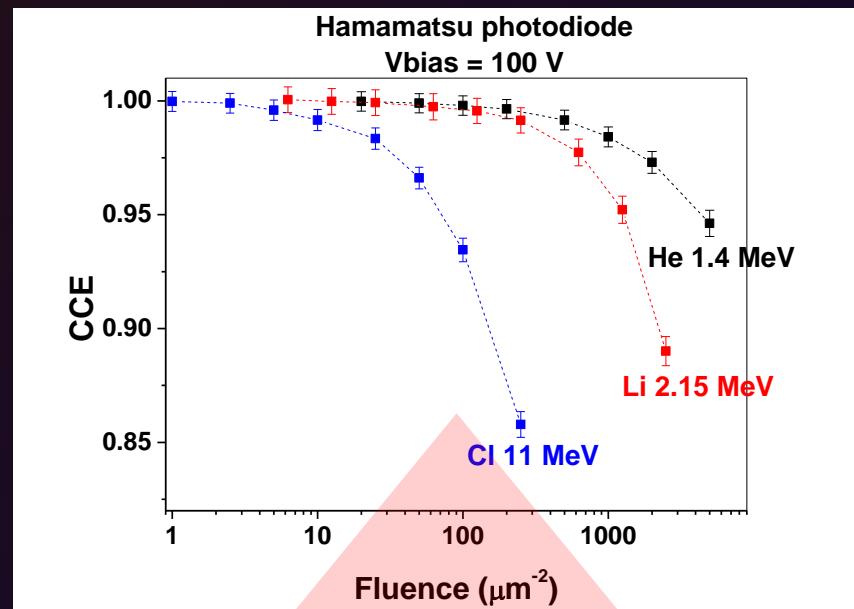


CCE
DEGRADATION



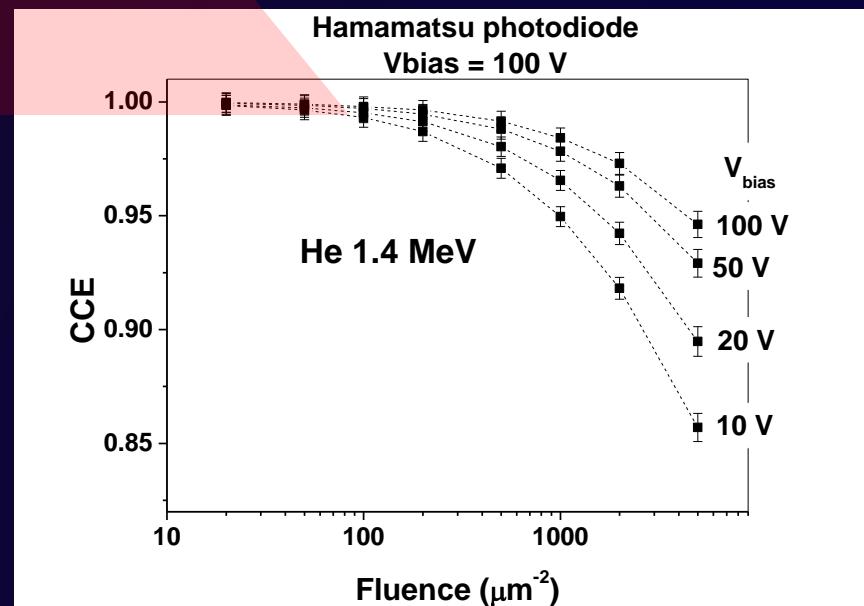


Damaging Ion Mass/Energy Fluence



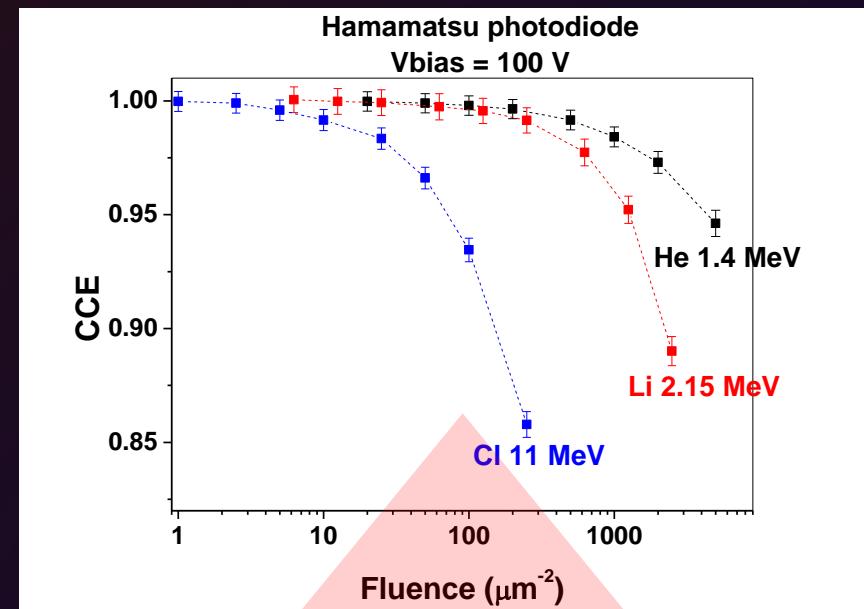
CCE
DEGRADATION

Electrostatics





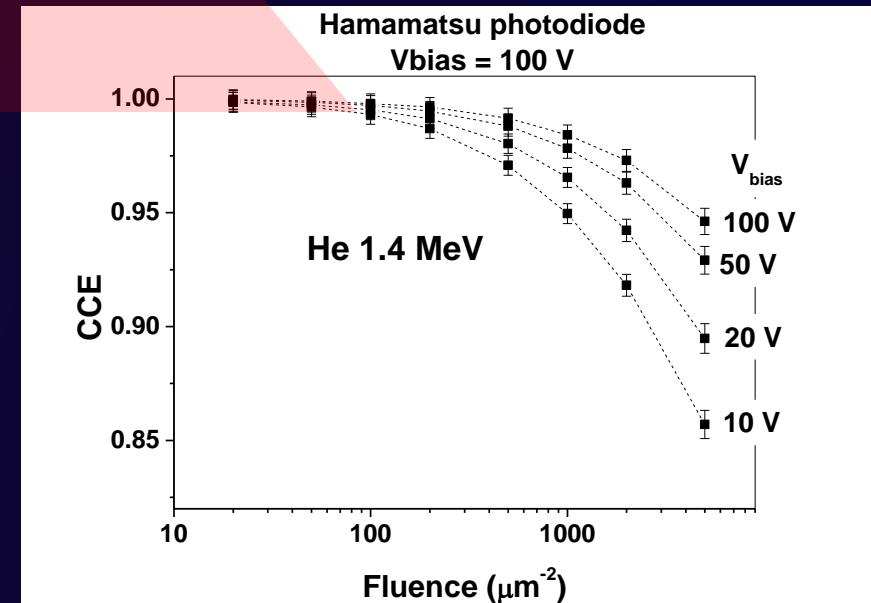
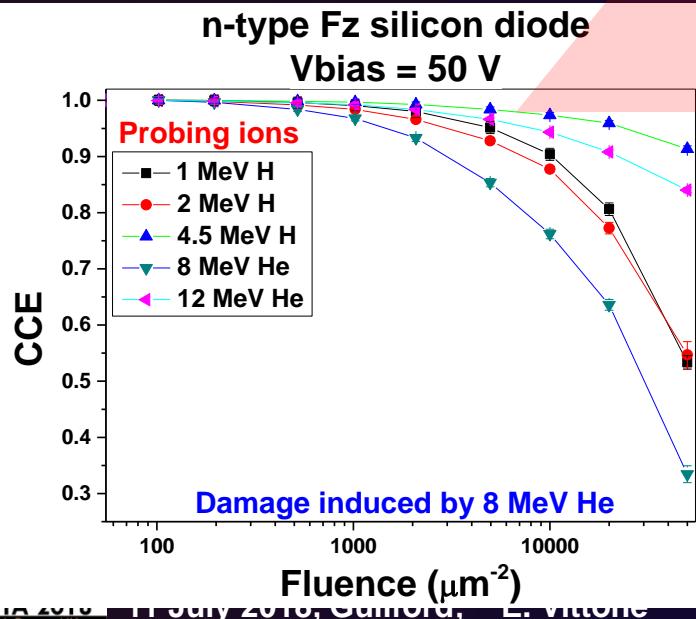
Damaging Ion Mass/Energy Fluence

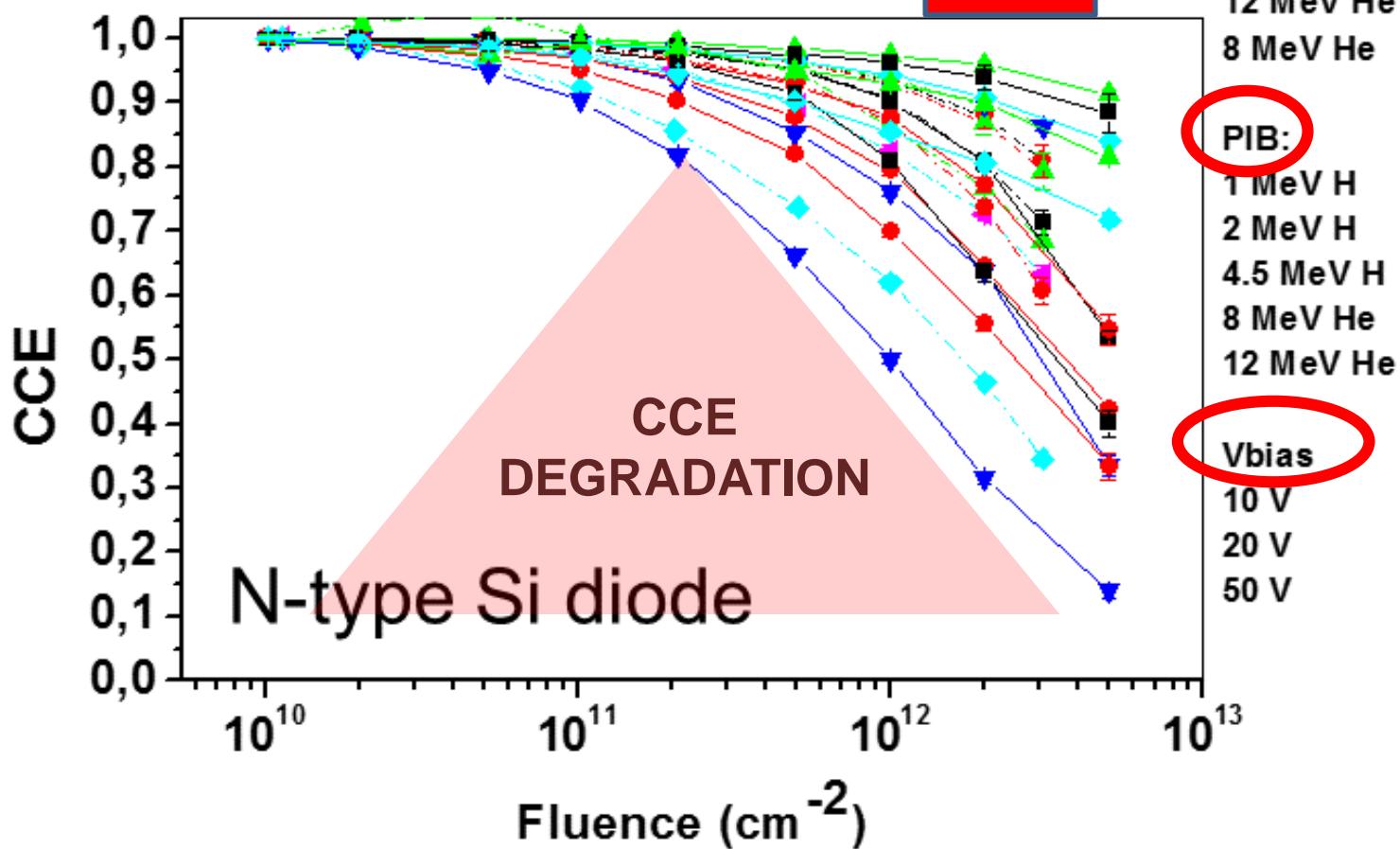


Probing Ion Mass/Energy

CCE DEGRADATION

Electrostatics







IAEA Coordinate Research Programme (CRP)

F11016 (2011-2015)



“Utilization of ion accelerators for studying and modeling of radiation induced defects in semiconductors and insulators”



CRP Outcome

A methodology to establish material parameters which reflect semiconductor radiation hardness by their ability to predict CCE degradation as a function of accumulated structural radiation damage.





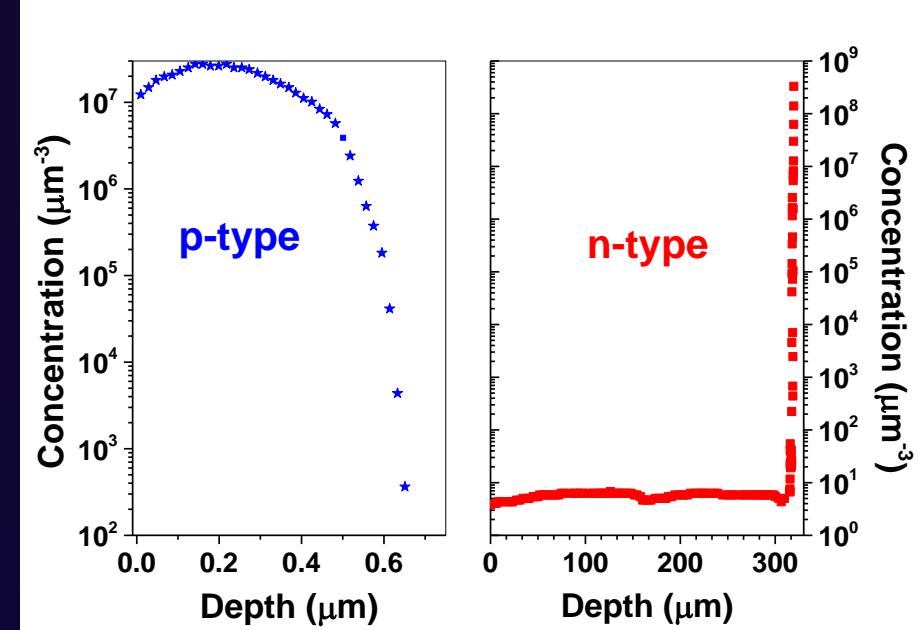
Sample under study

Commercially available p-i-n photodiode



Electrical characterization

Doping profile:
Spreding Resistance



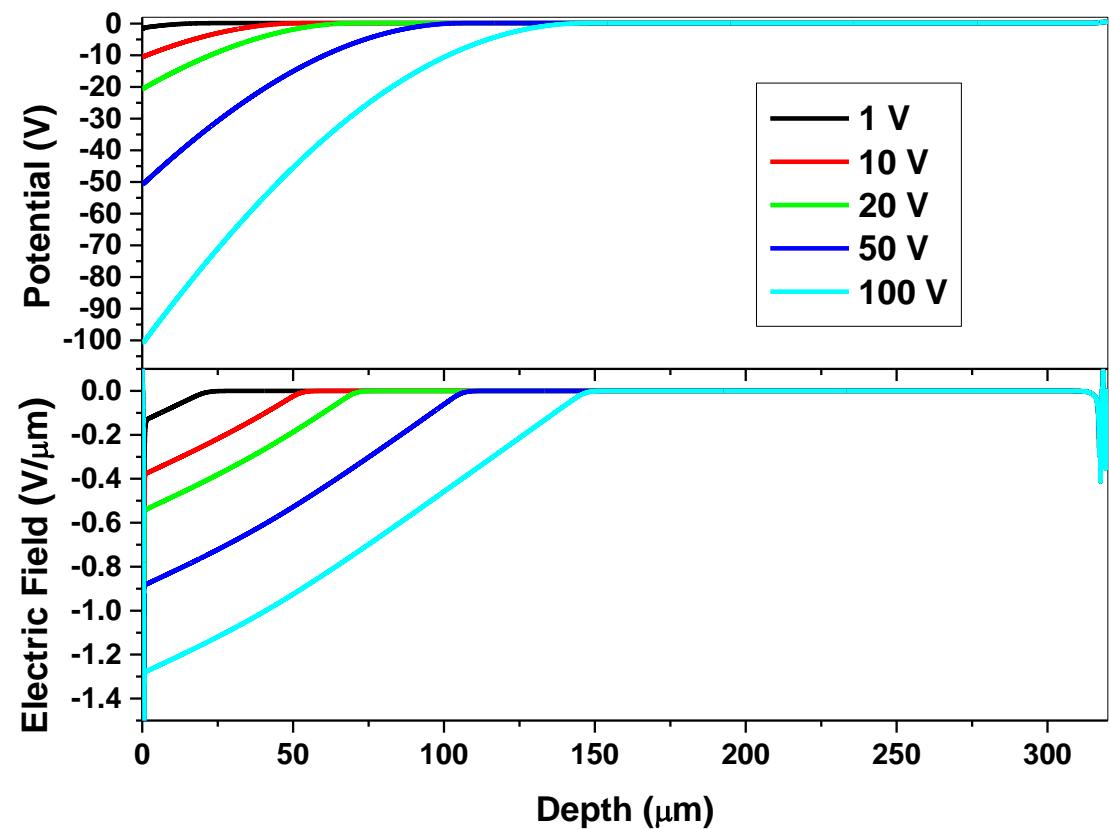


Sample under study

Commercially available p-i-n photodiode



Device Modeling
Electrostatics



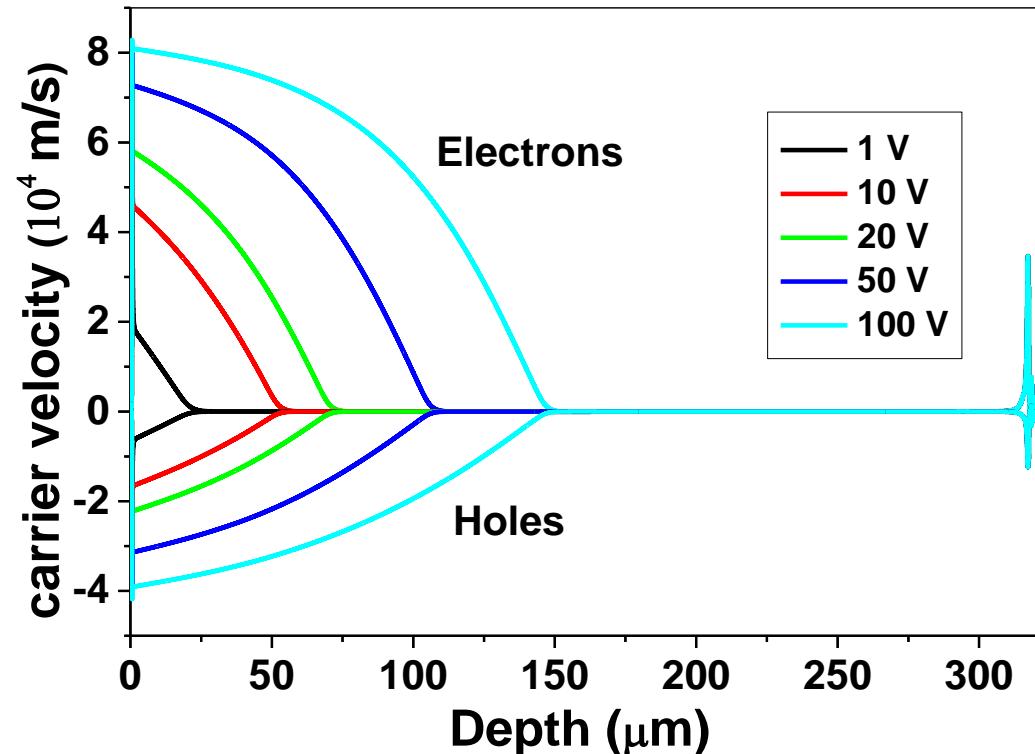


Sample under study

Commercially available p-i-n photodiode



Device Modeling
Transport



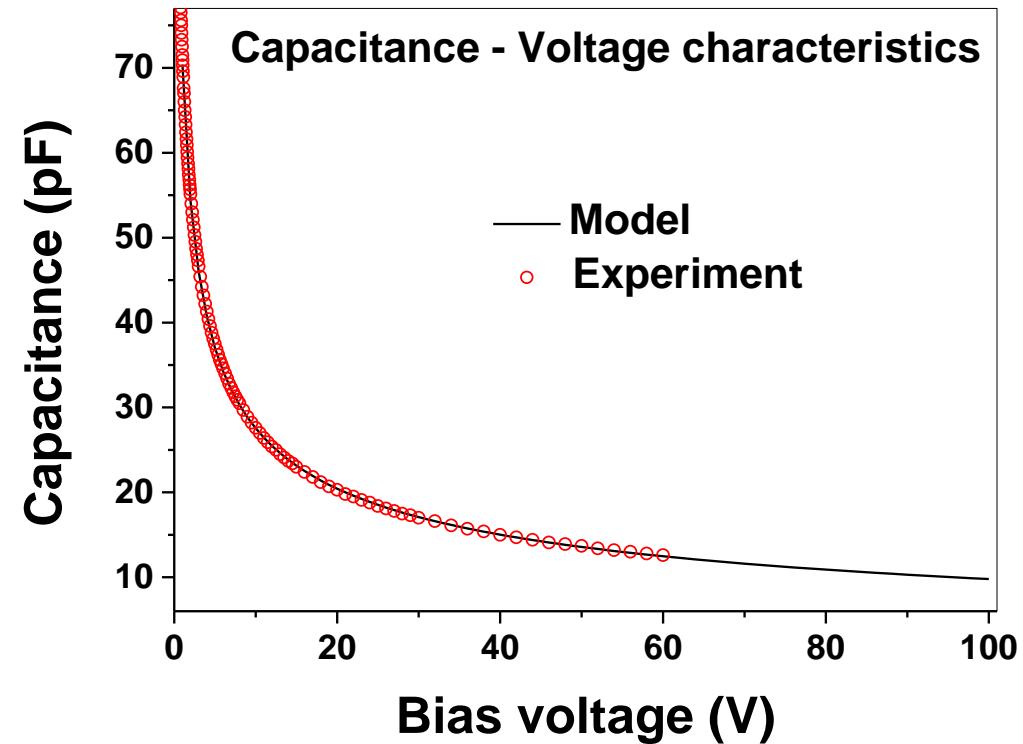


Sample under study

Commercially available p-i-n photodiode



Device Modeling Validation



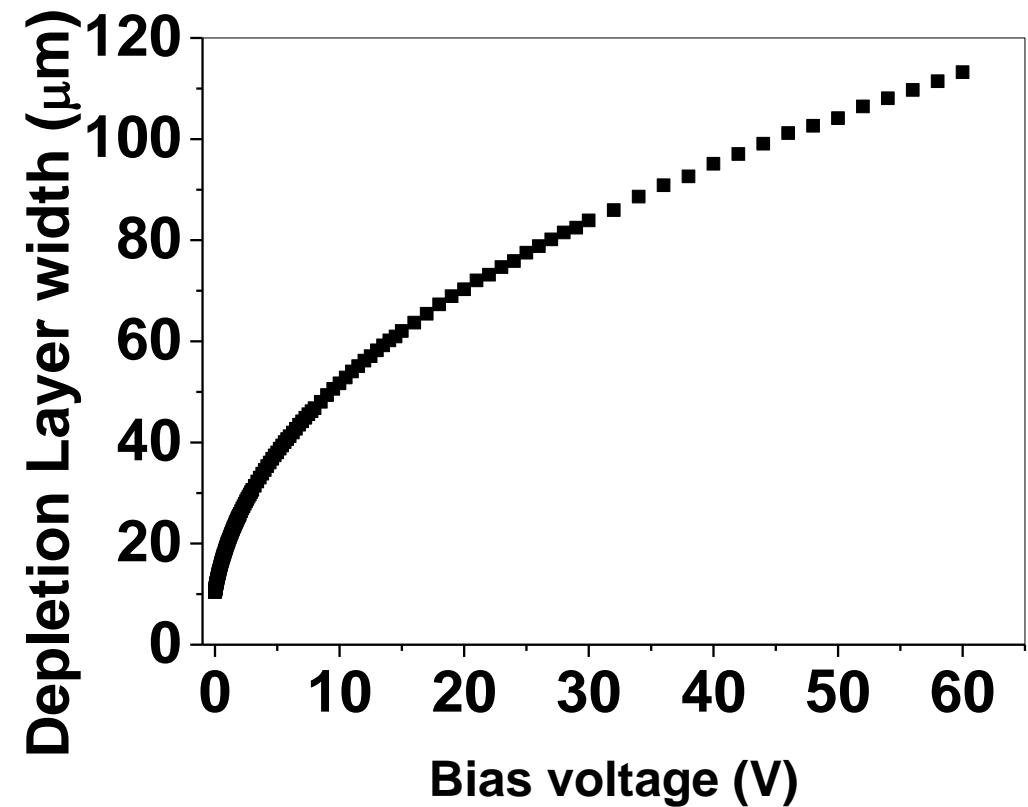


Sample under study

Commercially available p-i-n photodiode



Device Modeling Depletion Layer width



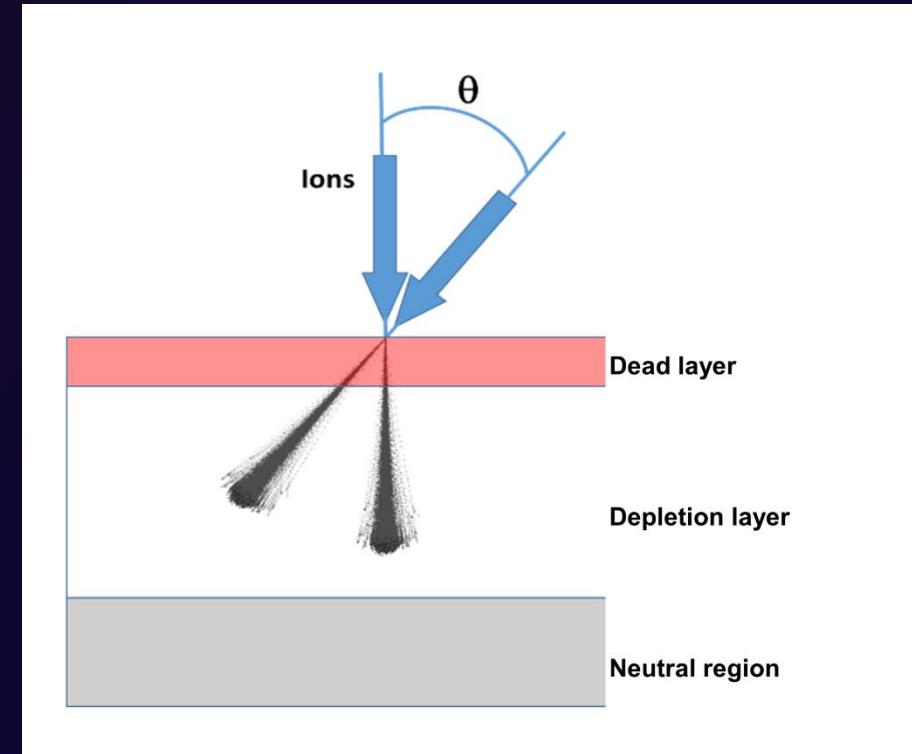


Sample under study

Commercially available p-i-n photodiode



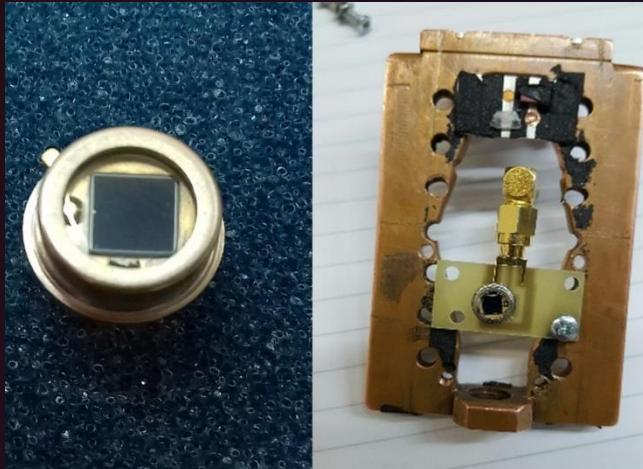
Dead layer
ARIBIC





Sample under study

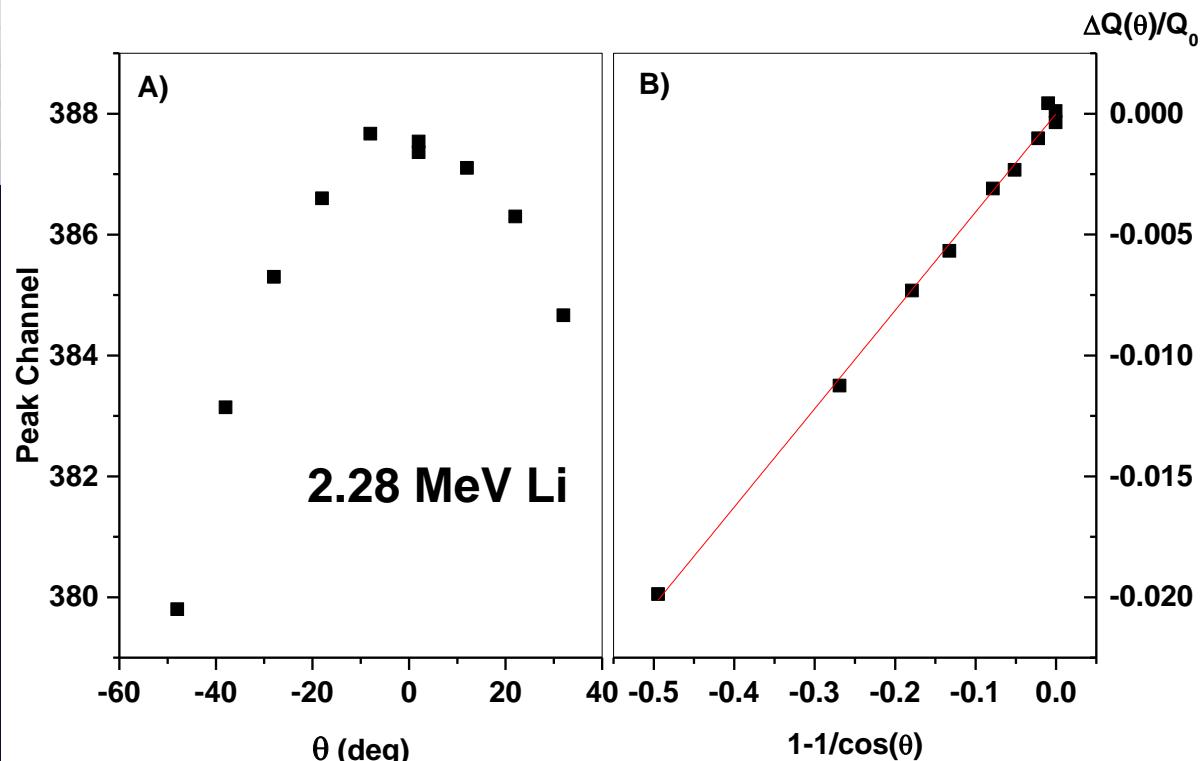
Commercially available p-i-n photodiode



Effective thickness in Si
 $t^* = 180 \text{ nm}$
RBS = 110 nm of SiO_2

Dead layer
ARIBIC

$$\frac{\Delta Q^*(\theta)}{Q_0} = \frac{1}{E_{\text{ion}}} \frac{dE}{dx} \Big|_0 t^* \cdot \left(1 - \frac{1}{\cos(\theta)}\right)$$



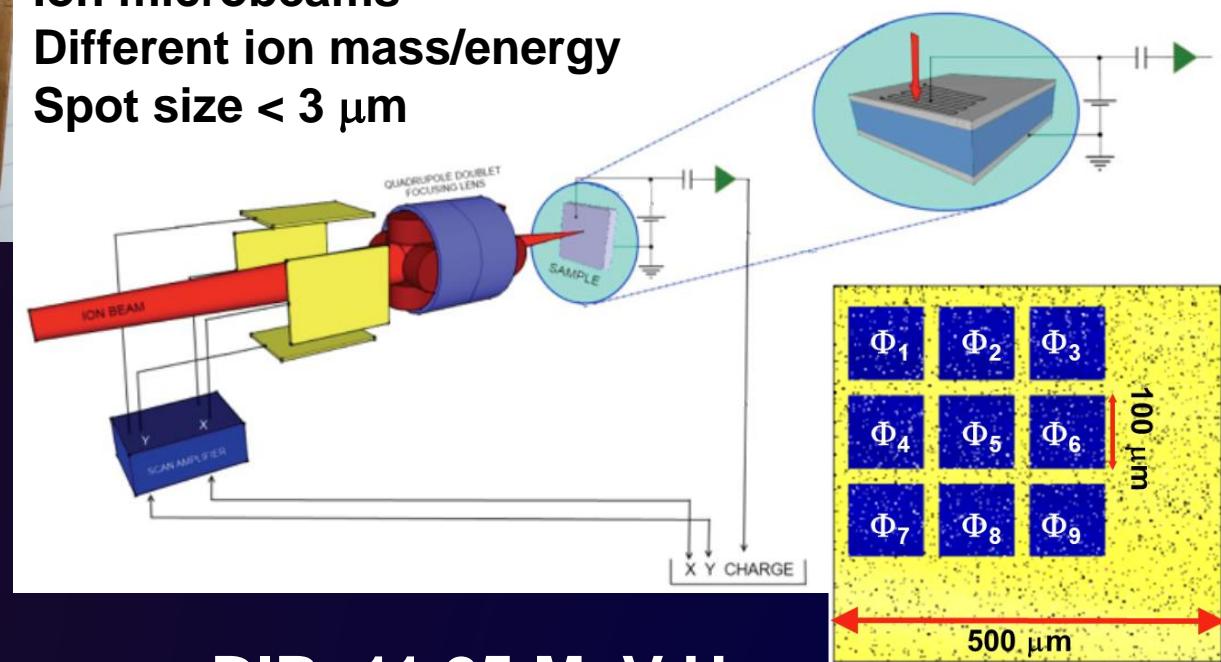
Sample under study

Commercially available p-i-n photodiode



Inducing the damage

Ion microbeams
Different ion mass/energy
Spot size < 3 μm



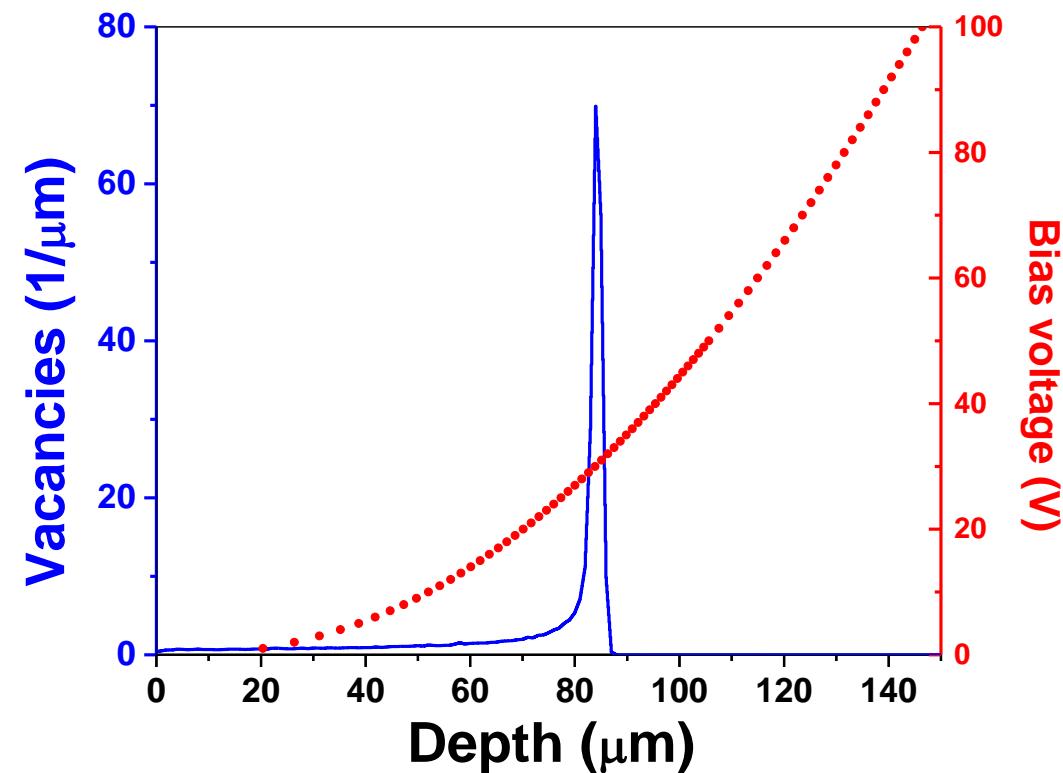


Sample under study

Commercially available p-i-n photodiode



Inducing the damage



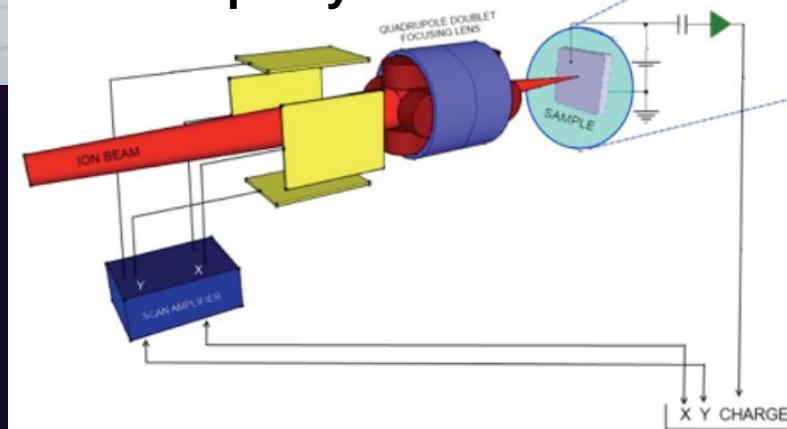
Sample under study

Commercially available p-i-n photodiode



Probing the damage

Ion microbeam
Low current (<fA)
CCE maps by IBIC



PIB
1.40 MeV He
11.25 MeV He
2.30 MeV H

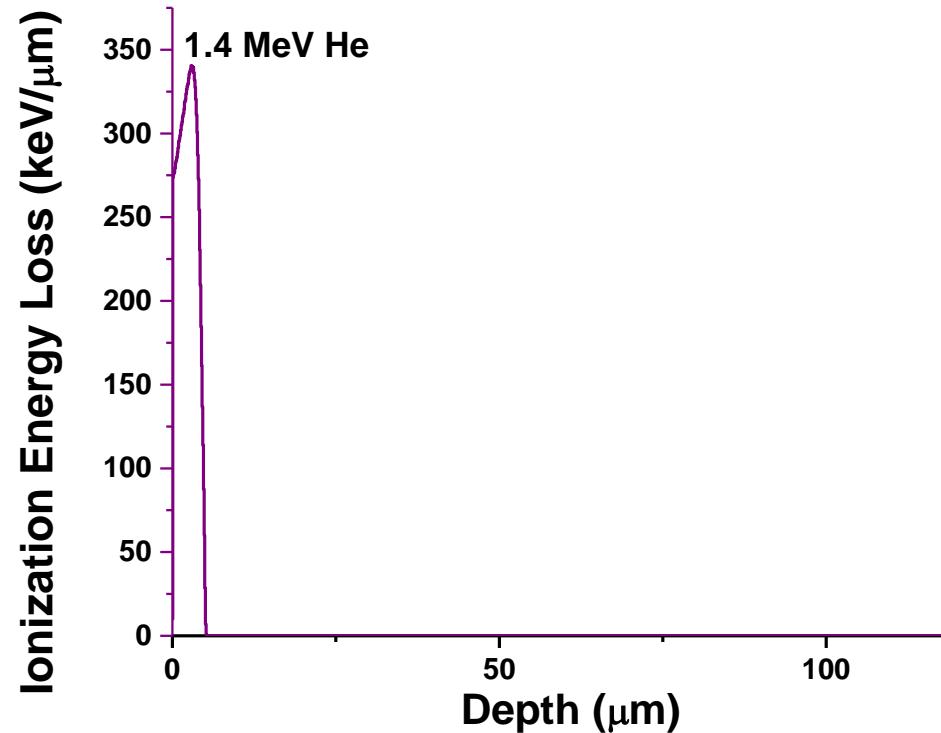


Sample under study

Commercially available p-i-n photodiode



Probing the damage



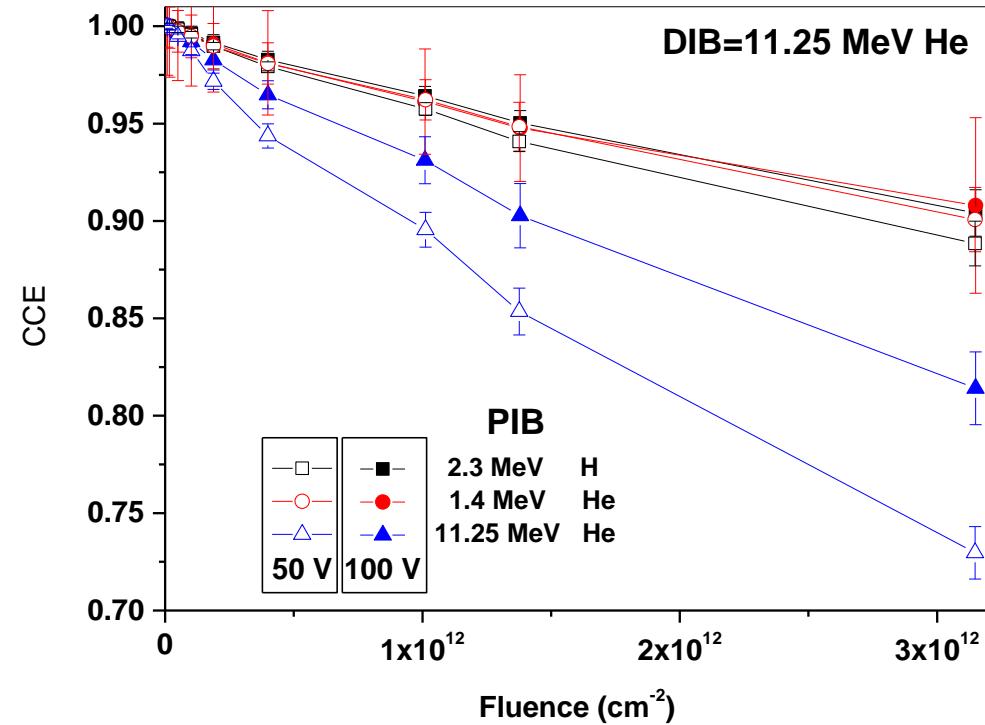


Sample under study

Commercially available p-i-n photodiode



Probing the damage
CCE at different bias
from different PIBs





Model for charge pulse formation (IBIC theory)

based on the Shockley-Ramo-Gunn theorem

Model for CCE degradation Based on the Shockley-Read-Hall model

Basic assumption:

- 1) In the linear regime, the ion induced damage affects mainly the carrier lifetime τ
- 2) The ion induced trap density is proportional to the VACANCY DENSITY

$$\frac{1}{\tau} = \frac{1}{\tau_0} + \alpha \cdot \text{Vac}(x) \cdot \Phi$$

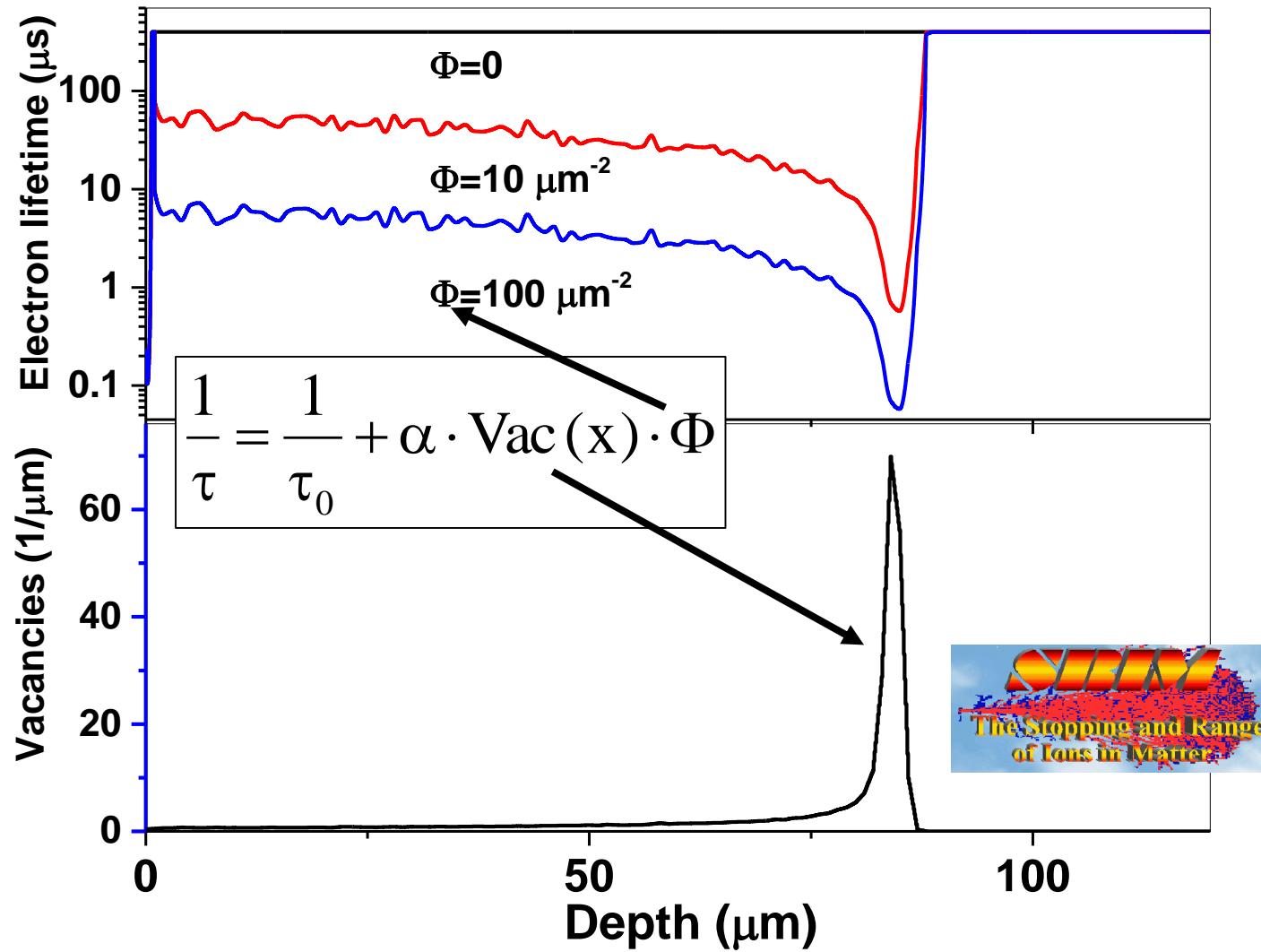
Diagram illustrating the model:
- A red box highlights the term $\alpha \cdot \text{Vac}(x) \cdot \Phi$.
- A green arrow labeled "Capture coefficient" points to the term α .
- A green arrow labeled "Vacancy Density Profile" points to the term $\text{Vac}(x)$.
- A green arrow labeled "Fluence" points to the term Φ .



Low level of damage: $\Phi < 10^{12} \text{ cm}^{-2} = (100 \times 100) \text{ nm}^2$

LOW DENSITY OF TRAPS -> NOT INTERACTING TRAPS







Model for charge pulse formation (IBIC theory)

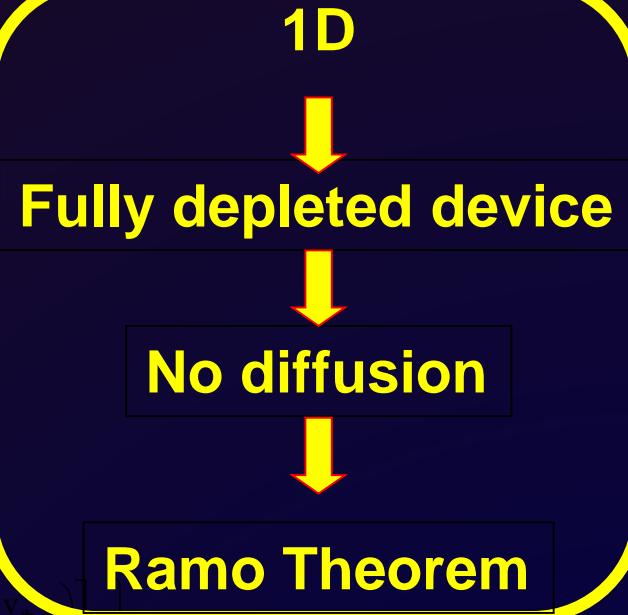


Ionization profile



Gunn's weighting field

$$Q_s = q \cdot \int_0^d dx \cdot \Gamma(x) \left\{ \int_x^d dy \cdot \frac{\partial F(y)}{\partial V_s} \cdot \exp \left[- \int_x^y dz \left(\frac{1}{v_p \cdot \tau_p} \right) \right] + \int_0^x dy \cdot \frac{\partial F(y)}{\partial V_s} \cdot \exp \left[- \int_y^x dz \left(\frac{1}{v_n \cdot \tau_n} \right) \right] \right\}$$

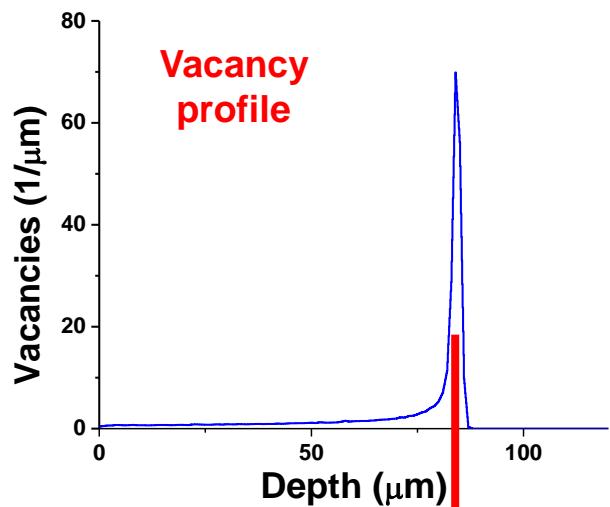


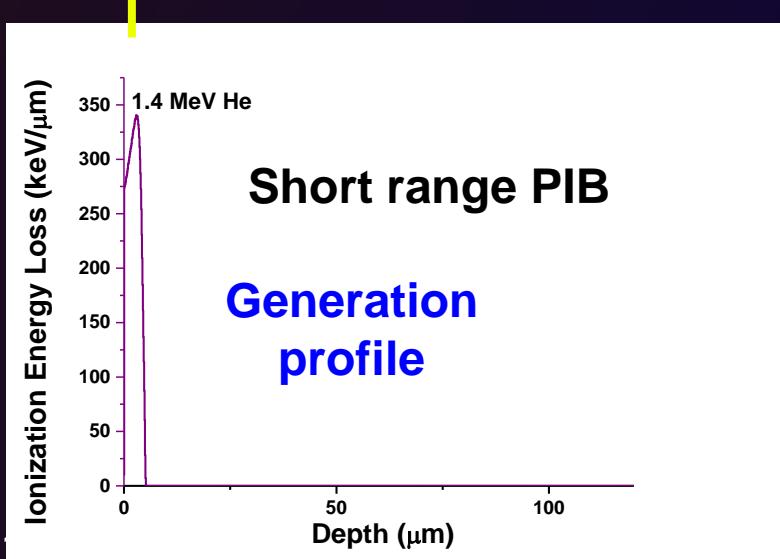
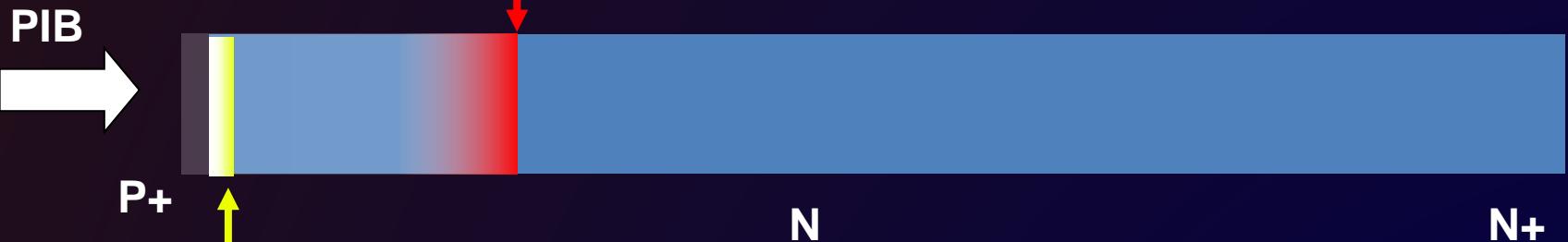
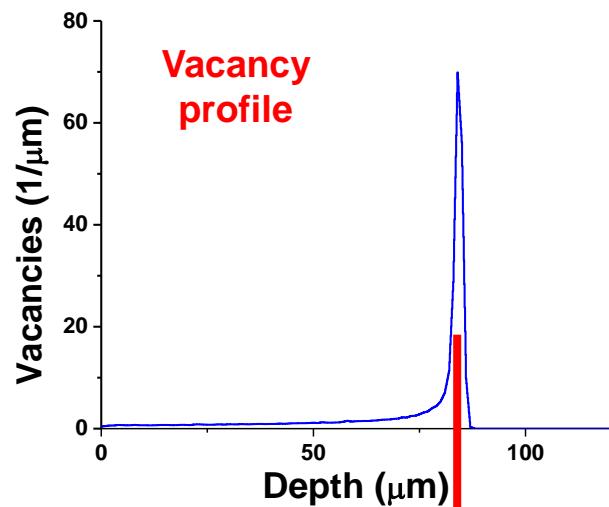
Holes

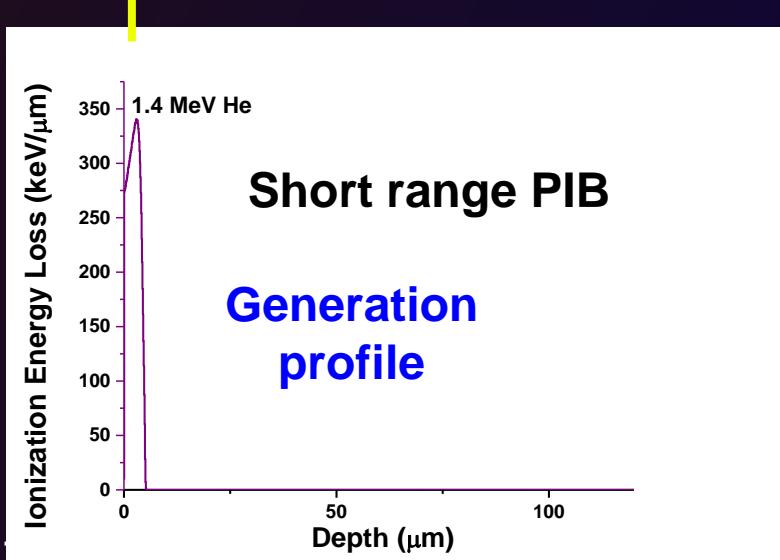
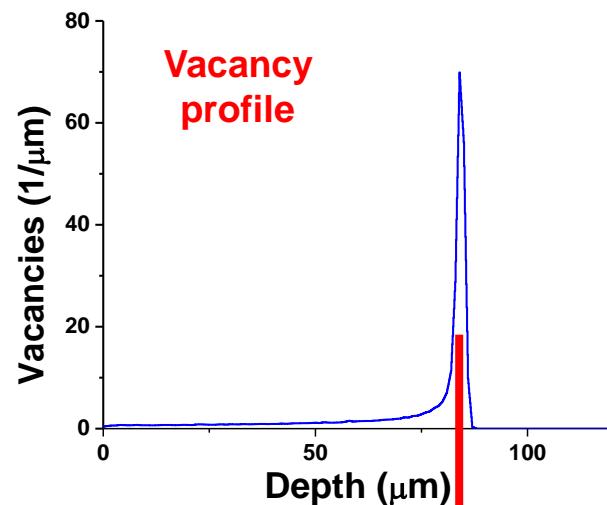
Electrons

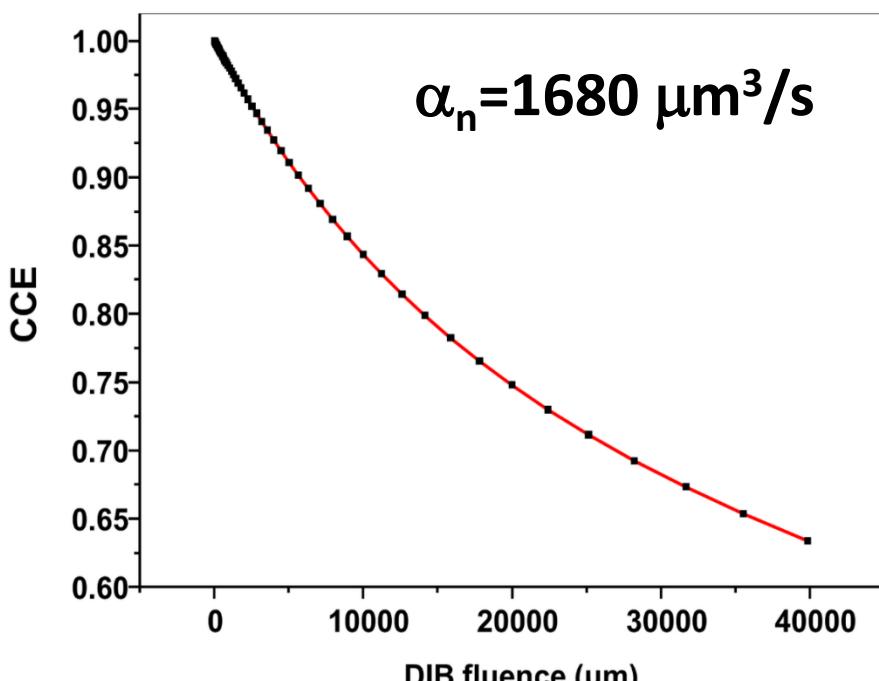
Drift lengths



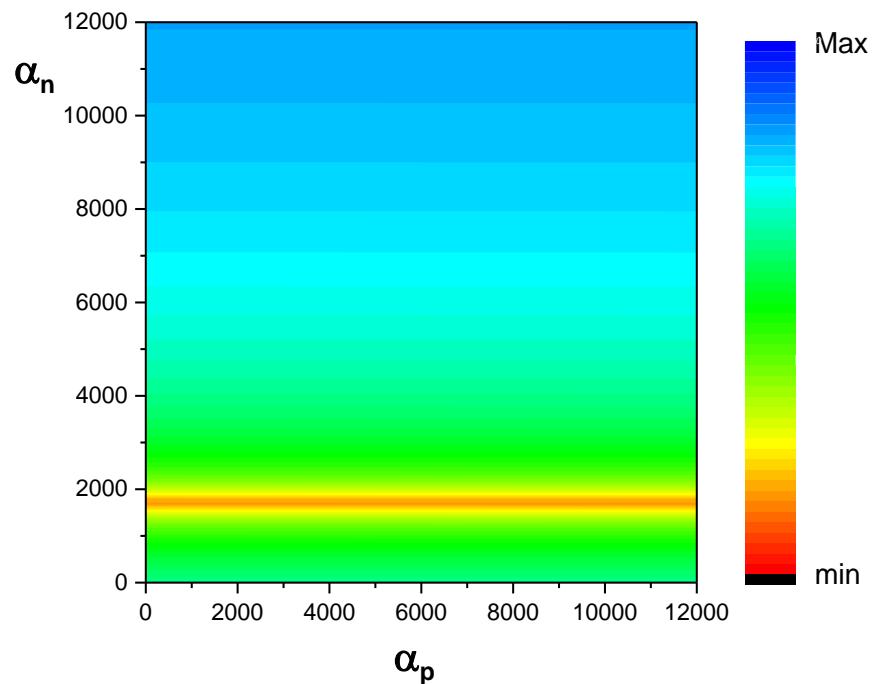






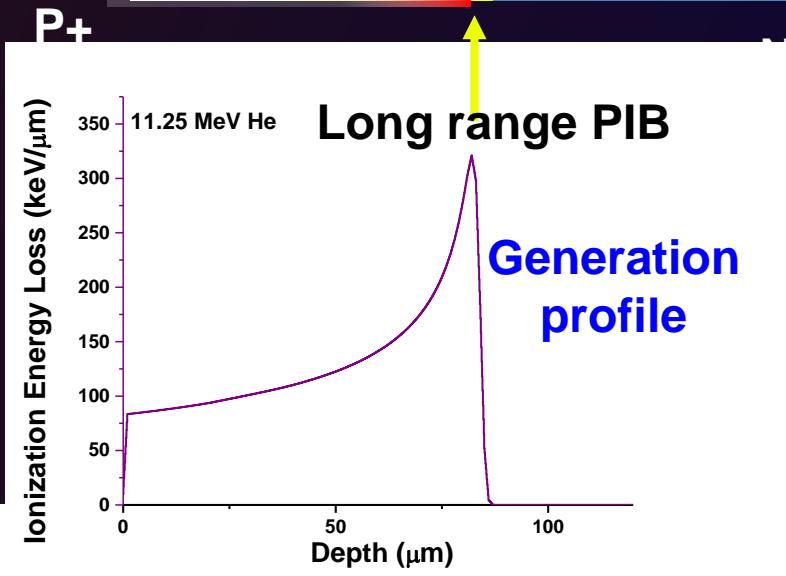
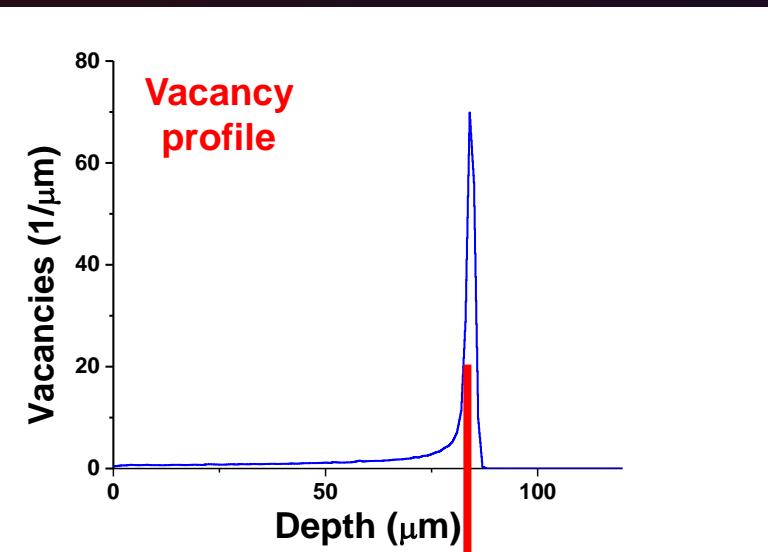


Residual map

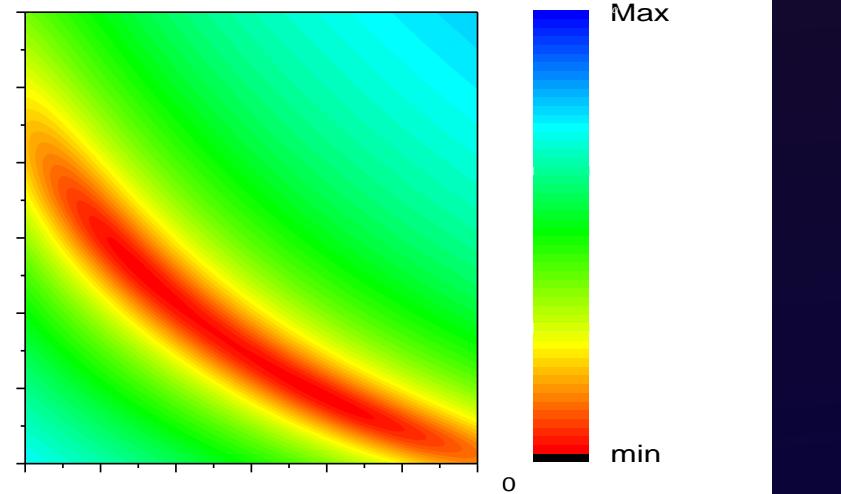
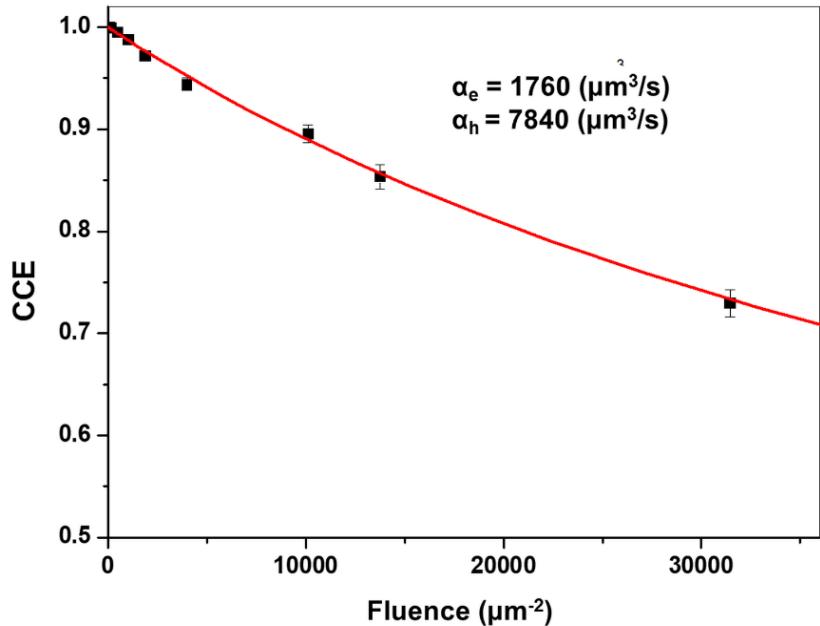


α_n Recombination Coefficient Free parameter

$$Q_s = q \cdot \int_0^d dx \cdot \Gamma(x) \left\{ \int_x^d dy \cdot \frac{\partial F(y)}{\partial V_s} \cdot \exp \left[- \int_x^y dz \frac{1}{V_n} \cdot \left(\frac{1}{\tau_0} + \alpha_n \cdot \text{Vac}(x) \cdot \Phi \right) \right] \right\}$$



Residual map

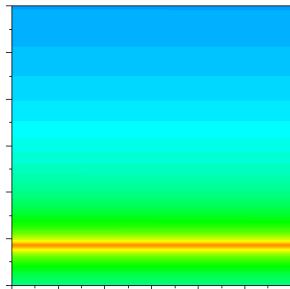


α_n, α_p
Recombination Coefficients
Free parameters

$$Q_s = q \cdot \int_0^d dx \cdot \Gamma(x) \left\{ \begin{aligned} & \left[\int_0^x dy \cdot \frac{\partial F(y)}{\partial V_s} \cdot \exp \left[- \int_y^x dz \frac{1}{V_p} \cdot \left(\frac{1}{\tau_0} + \alpha_p \cdot \text{Vac}(x) \cdot \Phi \right) \right] \right] + \\ & \left[\int_x^d dy \cdot \frac{\partial F(y)}{\partial V_s} \cdot \exp \left[- \int_x^y dz \frac{1}{V_n} \cdot \left(\frac{1}{\tau_0} + \alpha_n \cdot \text{Vac}(x) \cdot \Phi \right) \right] \right] \end{aligned} \right\}$$



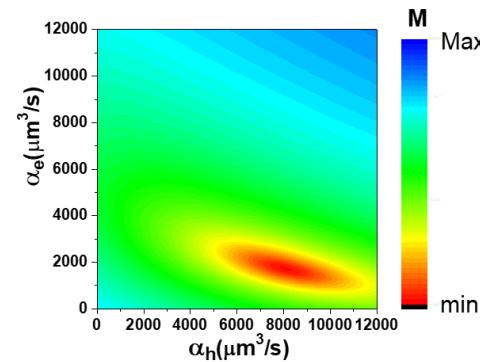
PIB=1.4 MeV He



Short range PIB

+

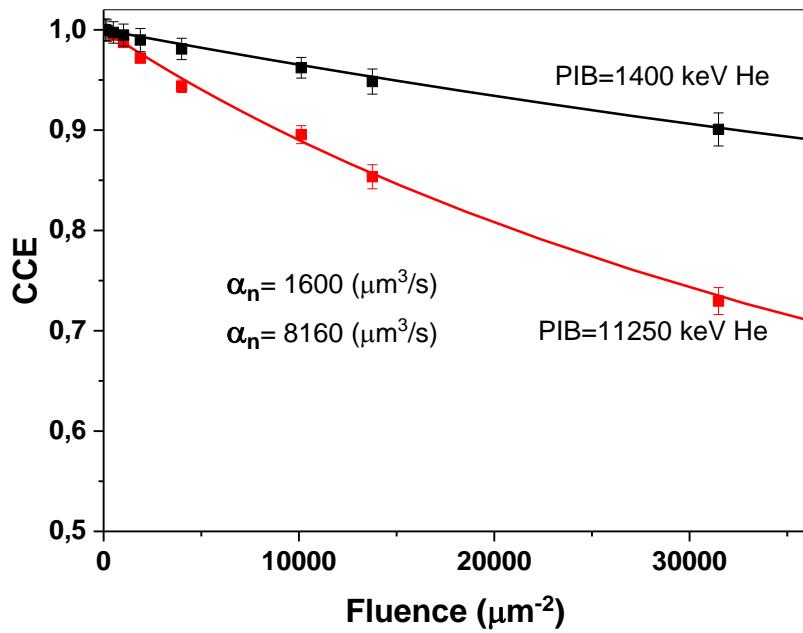
=



M Max
min

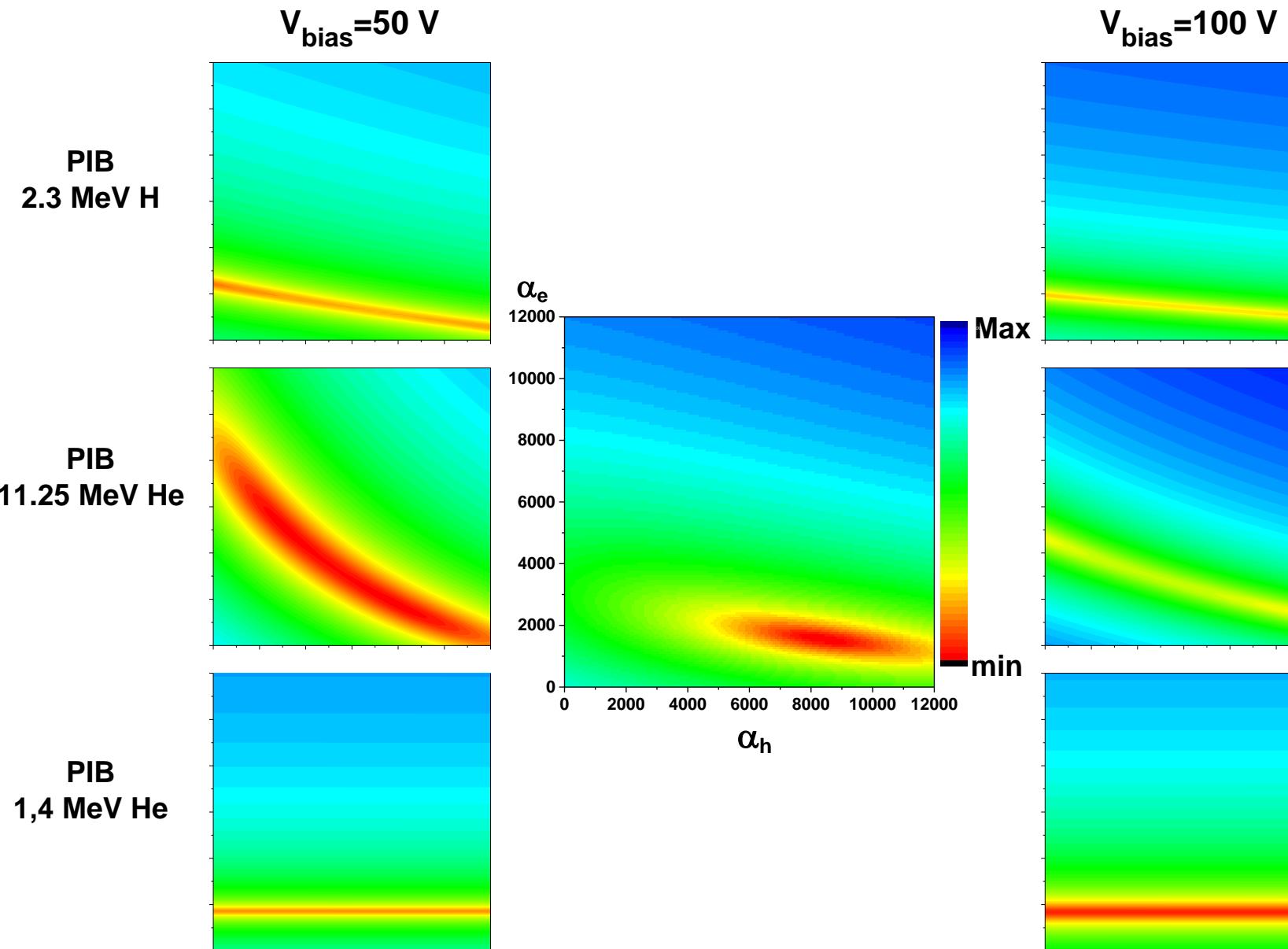
Long range PIB

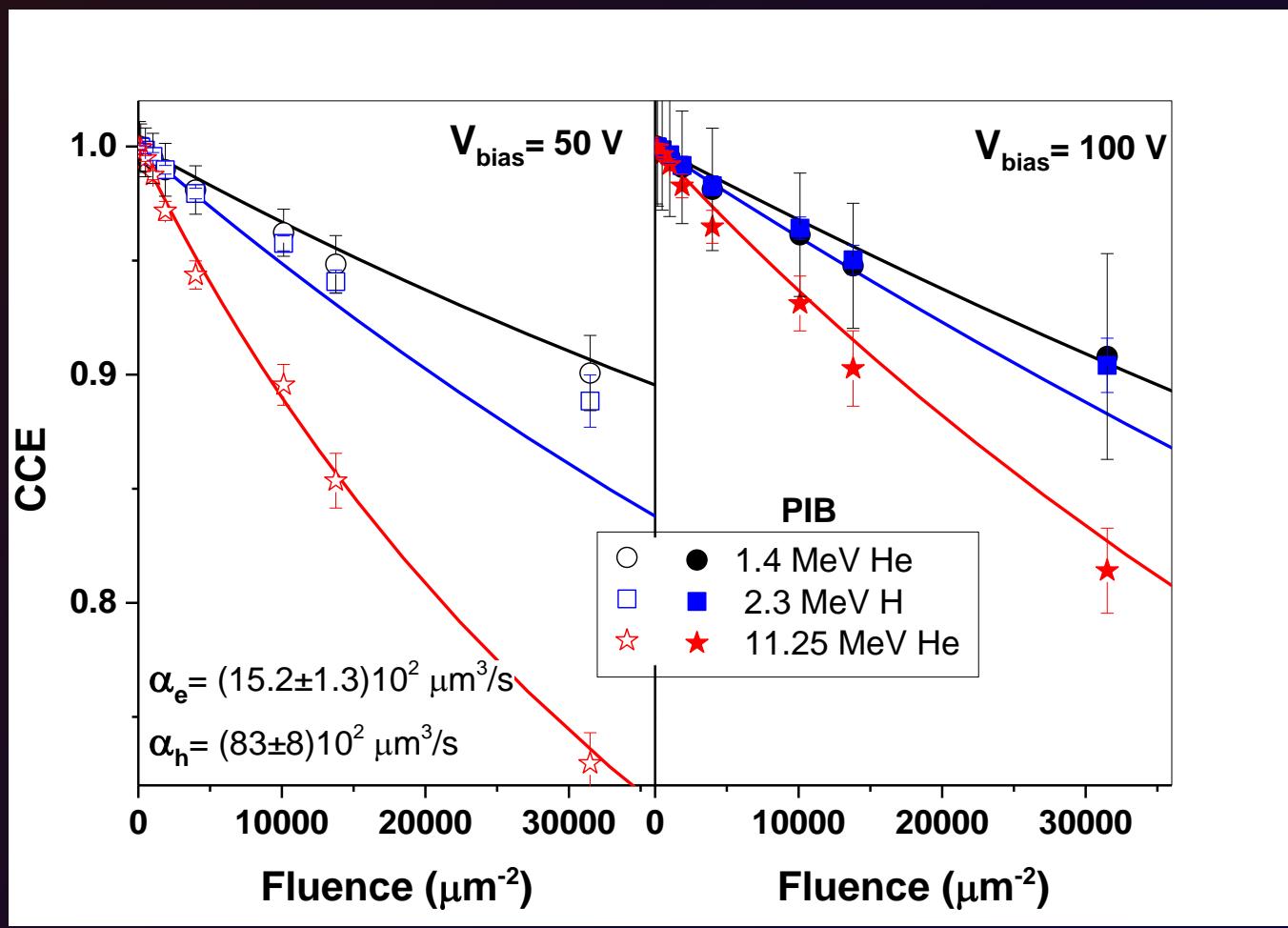
PIB=11.25 MeV He



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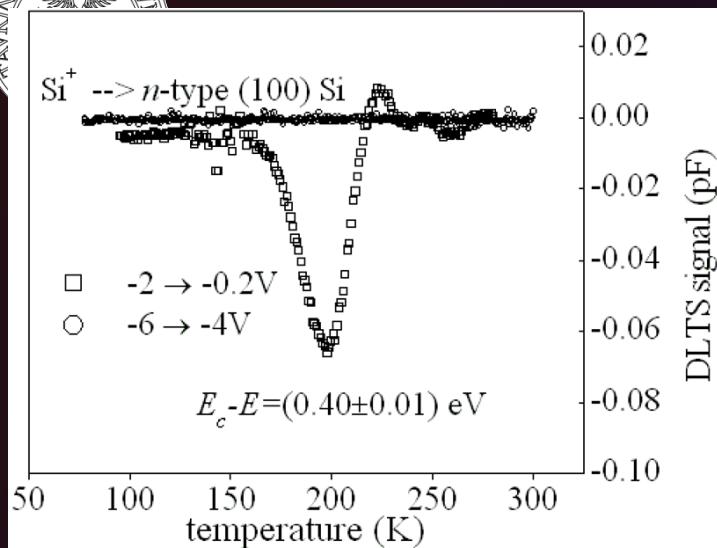
11 July 2018, Guilford; E. Vittone





The solid lines are the best fits obtained by means of our model considering

- Different PIBs
- Different biases (50 V, 100 V)



N-type silicon
DLTS measurements
singly V2(-/-0) negatively charged
divacancy

$$\sigma_n \approx 5 \cdot 10^{-15} \text{ cm}^2$$

$$\alpha_n \approx 1520 \cdot 10^{-12} \text{ cm}^3/\text{s}$$

$$v_{th} \approx 2.05 \cdot 10^7 \text{ cm/s}$$

$$\alpha = k \cdot \sigma \cdot v_{th} \longrightarrow k = 0.016$$



about 60 radiation induced vacancies are required to form one stable electron recombination centre.



Limits of applicability

Basic Hypotheses

DIB : low level of damage

$$\frac{1}{\tau_{e,h}} = \frac{1}{\tau_{0,e,h}} + \alpha_{n,p} \cdot \text{Vac}(x) \cdot \Phi = \frac{1}{\tau_{0,e,h}} + (\sigma_{e,h} \cdot v_{th}) \cdot \text{Vac}(x) \cdot \Phi$$

“linear model”
Independent traps, no clusters

Unperturbed electrostatics (i.e. doping profile) of the device

PIB : ion probe

CCE is the sum of the individual e/h contributions

No plasma effects induced by probing ions





Recombination coefficient: $\alpha = k \cdot \sigma \cdot v_{th}$



| Ref. | Diode | PIBs | DIBs | Max Fluence (μm^{-2}) | α_e ($\mu\text{m}^3/\text{s}$) | α_h ($\mu\text{m}^3/\text{s}$) |
|------|--|---|---|---------------------------------------|--|--|
| [2] | Hamamatsu S5821 | 1.4 MeV He | 1.4 MeV He 2.15 MeV Li 4.0 MeV O 11.0 MeV Cl | 5000 2000 500 200 | 8800 ± 1200 | -- |
| [3] | Hamamatsu S5821* | 1.036 MeV He | 1.036 MeV He 2 MeV He | 4000 | 10270 ± 260 | 23500 ± 2800 |
| [1] | n.type Si PIN diode from Helsinki University | 2 MeV He 2 MeV H 8 MeV He 12 MeV He 4.5 MeV H | 4 MeV He 8 MeV He | 20000 | 2500 ± 300 | 210 ± 160 |
| [1] | p.type Si PIN diode from Helsinki University | 2 MeV He 2 MeV H 8 MeV He 12 MeV He 4.5 MeV H | 4 MeV He 8 MeV He | 20000 | 2200 ± 300 | 1310 ± 90 |
| [4] | Hamamatsu S1223 | 1.4 MeV He 2.3 MeV H 11.25 MeV He | 11.25 MeV He | 30000 | 1520 ± 130 | 8300 ± 800 |

[1] E. Vittone et al. Nuclear Instr. and Methods in Physics Research B 372 (2016) 128–142

[2] Ž. Pastuović et al. Applied Physics Letters 98, 092101 (2011)

[3] J. Garcia et al. Unpublished

[4] This work



CONCLUSIONS



The IAEA methodology has been used to study the radiation hardness of a commercially available silicon p-i-n diode. This methodology contribute towards a standardized quantification of radiation hardness of semiconductor materials.

The capture coefficient is directly related to the radiation hardness of the material

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